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THE ANISOTROPY OF THE ELECTRICAL RESISTIVITY AND THERMAL EXPANSION OF SINGLE CRYSTAL DyNi_5

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The results of the measurements of temperature dependence of the electrical resistivity $\rho(T)$, thermal expansion $L(T)$ and susceptibility $\chi(T)$ on single crystal DyNi_5 in the temperature range 4.2–300 K are reported. The maximum in $\rho(T)$ at $T_N = 12.2$ K was observed which corresponds to the transition to antiferromagnetically ordered state. Taking into account the Bloch-Grüneisen formula for electron-phonon contribution $\rho_{\text{ph}}(T)$ to $\rho(T)$ the anisotropic magnetic contribution $\rho_{\text{mag}}(T)$ was determined. We also observed the anisotropic behaviour of $L(T)$ at low temperatures.

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The intermetallic RNi_5 (R = rare earths) compounds crystallize in the simple hexagonal structure and have been the subject of detailed studies during the last years [1]. Their magnetic properties are quite well understood and they are predominantly determined by the crystalline electric field (CEF) acting on R and the exchange interactions. In this series the Ni atoms are non-magnetic, although they are very close to the onset of Ni magnetism. However, till now there is only a few pieces of information about the anisotropy of the electrical transport in LaNi_5 and PrNi_5 [2] and only the information about the thermal expansion of PrNi_5 [3]. From the heat capacity measurements follows that the intermetallic compound DyNi_5 undergoes an antiferromagnetic ordering transition at $T_N = 11.6$ K [4]. Most of information about CEF level scheme of DyNi_5 has been obtained indirectly from specific heat [4], magnetization and susceptibility [5, 6] in the paramagnetic state. The electron-quasiparticle interaction function was measured by point-contact (PC) spectroscopy above and below T_N [7]. In order to study the anisotropy of the CEF contribution to the transport and thermal properties, for the first time we performed the measurements of the electrical resistivity and the

thermal expansion coefficient α on the oriented single crystal of DyNi₅. The susceptibility measurements were performed too.

The single crystal of a good quality was grown by the vertical floating zone method in an image furnace. The standard four probe d.c. method was used for the electrical resistivity measurements in the temperature range 4.2–300 K along 3 main crystallographical axis. The d.c. susceptibility measurements were performed in the vibrating sample magnetometer. The thermal expansion measurements were performed by resistive strain gauge technique within temperature range 4.2–300 K perpendicular and parallel to basal plane.

The temperature dependences of the electrical resistivity $\rho(T)$ of DyNi₅ measured along the main crystallographical axis are shown in Fig. 1a. One could see the typical metallic behaviour. Taking into account the residual resistivities ρ_0 ($\rho_0^a = 9.9 \mu\Omega \text{ cm}$, $\rho_0^b = 11.7 \mu\Omega \text{ cm}$ and $\rho_0^c = 8.5 \mu\Omega \text{ cm}$), the following residual resistivity ratios $\text{RRR} = \rho(300 \text{ K})/\rho_0$ have been obtained: $\text{RRR}_a = 4.4$, $\text{RRR}_b = 4.0$, $\text{RRR}_c = 3.9$. The low temperature behaviour of the electrical resistivity ($\rho(T) - \rho_0$) above and below T_N is presented in Fig. 1b. The clear maximum at $T_N = 12.2 \text{ K}$ is connected with the transition to the antiferromagnetic state. This temperature is in a good agreement with previous data [4–6].

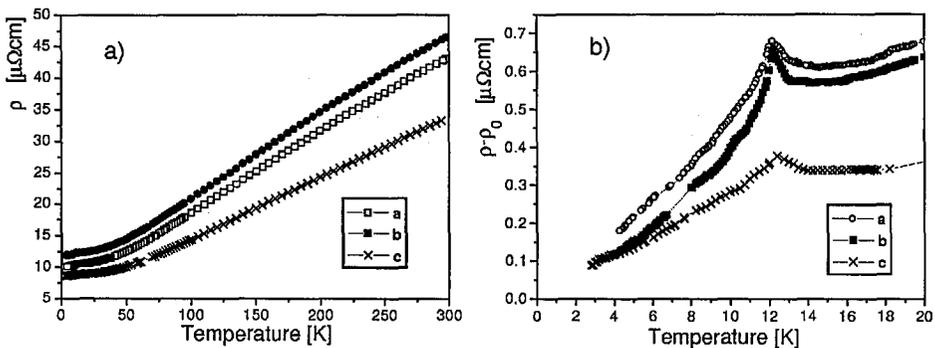


Fig. 1. The temperature dependences of the electrical resistivity of single crystal DyNi₅ along main crystallographical axes (a). The low temperature detail is presented in (b).

In order to determine the magnetic contribution to the electrical resistivity which is connected with CEF we subtracted the residual resistivity and phononic contribution in a similar manner as in [2]. We used for the subtraction of phononic term the Bloch–Grüneisen formula for electron–phonon scattering $\rho_{\text{ph}}(T)$. We obtained the best results with the Debye temperature $\Theta_D = 380 \text{ K}$. The magnetic contribution $\rho_{\text{mag}}(T)$ (Fig. 2) is anisotropic and it follows the characteristic shape [2] with high temperature saturation ($\rho_{\text{mag}}^a(300 \text{ K}) = 4.5 \mu\Omega \text{ cm}$, $\rho_{\text{mag}}^b(300 \text{ K}) = 4.9 \mu\Omega \text{ cm}$ and $\rho_{\text{mag}}^c(300 \text{ K}) = 1.8 \mu\Omega \text{ cm}$). All dependences contain the maximum at T_N . This behaviour is characteristic of the scattering of the conduction electrons on the disordered magnetic moments (spin disorder resistivity model [2]) in a presence of the CEF. The anisotropy of the magnetic contribution to the electrical resistivity of DyNi₅ is similar to PrNi₅ [2] where

the contribution along c axis is about 2.5 times smaller than along b and a . Also $\rho_{\text{mag}}^b(T) \geq \rho_{\text{mag}}^a(T)$. Due to this behaviour we could expect the anisotropic CEF contribution to the electron-quasiparticle interaction function which has been measured by PC spectroscopy [7] along a axis. This contribution will be smaller along c axis than along b or a axis.

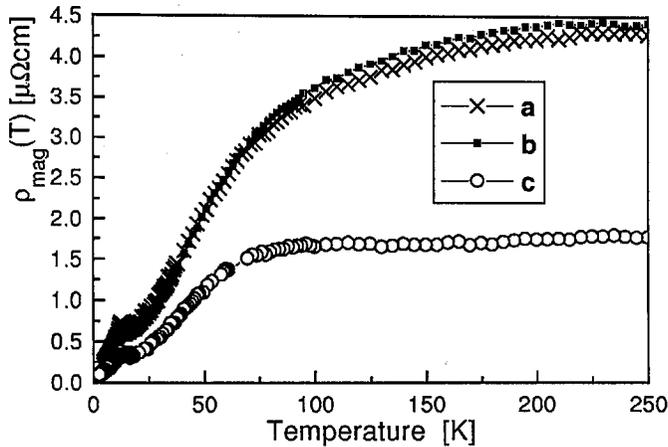


Fig. 2. The magnetic contributions $\rho_{\text{mag}}(T)$ to the the electrical resistivity of single crystal DyNi_5 along main crystallographical axes.

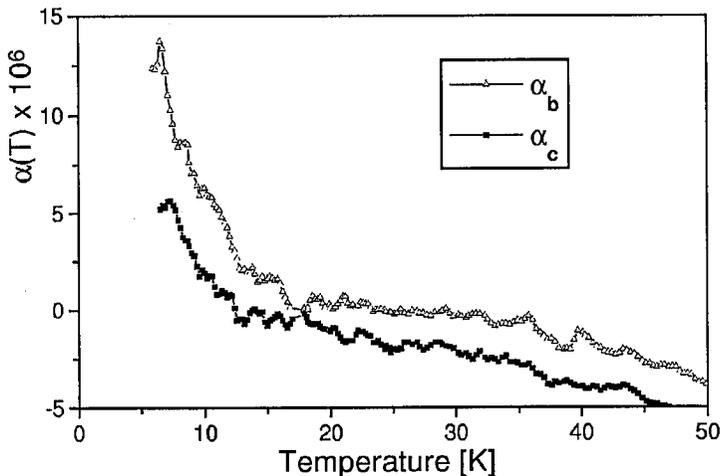


Fig. 3. The temperature dependences of the linear thermal expansions coefficients α perpendicular and parallel to the basal plane of single crystal DyNi_5 .

The measurements of the temperature dependence of susceptibility along main crystallographical axis were performed in order to check the transition temperature T_N determined from electrical resistivity. We obtained the same value of

T_N for both methods. The susceptibility data are similar to the previous data [5] and they are characteristic of antiferromagnetic transition. We observed that the c axis is the hard axis of magnetization. The magnitudes of susceptibility in magnetically ordered state along b are slightly higher than along a .

In Fig. 3 the low temperature part ($T < 30$ K) of the linear thermal expansion coefficient $\alpha(T)$ perpendicular ($\alpha_c(T)$) and parallel ($\alpha_b(T)$) to basal plane are shown. At low temperatures below T_N both dependences become positive. At high temperatures the $\alpha(T)$ is always negative for both orientation and $\alpha_b(T) > \alpha_c(T)$. Both dependences saturate to constant value at room temperature 300 K ($\alpha_b(T) = -21 \times 10^{-6} \text{ K}^{-1}$ and $\alpha_c(T) = -22 \times 10^{-6} \text{ K}^{-1}$). On the other hand, one could see the change of slope of both dependences at about $T_N = 12.2$ K. Generally, the magnitudes of absolute values of $\alpha(T)$ are smaller than in case of PrNi_5 [3].

The electrical resistivity, susceptibility and thermal expansion of single crystal DyNi_5 along main crystallographical axis were measured. We estimated the anisotropic magnetic contribution to the electrical resistivity and the linear thermal expansion coefficient.

Acknowledgments

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