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PHOTOEMISSION AND INVERSE PHOTOEMISSION STUDIES OF SiO₂

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Occupied and unoccupied electron states of amorphous silicon dioxide film supported on Si crystal are studied by using X-ray photoemission and, for the first time, X-ray inverse photoemission (X-ray bremsstrahlung isochromat method). A special care was undertaken to minimize decomposition of silicon oxide during X-ray bremsstrahlung measurements. The experimental spectra are compared with theoretical band structure calculations for amorphous SiO₂ from the literature and good overall agreement is found.

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1. Introduction

Amorphous SiO₂ (referred hereafter as a-SiO₂), being a wide-gap insulator, is of considerable technological and theoretical interest, therefore a thorough investigation of the electronic states in a-SiO₂ is desirable. Vacuum ultraviolet (VUV) and X-ray spectroscopies are very suitable techniques for studies of electronic structure of materials. The occupied part of electron states can be probed by VUV and X-ray photoemission methods giving more or less a total density of states (DOS) in material, whereas X-ray emission band spectra probe a partial DOS and local electronic structure e.g. Si K_{β} emission spectrum gives electronic structure of *p*-type valence electrons around Si ion.

Numerous experimental works on VUV and X-ray photoemission [e.g. 1–3] and X-ray emission spectra [e.g. 4–6] have contributed to a proper understanding of the valence-band structure of SiO₂, although much less is known about unoccupied electron bands.

There are a few reports of unoccupied electron states in SiO₂ by measuring X-ray absorption spectra at the Si *K* edge [7–9], Si *L*_{2,3} edge [10–12] and O *K*

edge [10, 13, 14]. In general, a synchrotron radiation source is recommended for X-ray absorption studies.

Unoccupied electron states can be also studied by VUV and X-ray bremsstrahlung isochromat method (named inverse photoemission) giving the total density of states weighted by transition probability [15–17]. Surprisingly little was done in this field for SiO_2 . To our knowledge, only two articles dealing with inverse photoemission studies of SiO_2 were published up to now [18, 19]. Both studies were performed in VUV photon range and for very thin oxide layers on an Si surface. This is partially due to general limitation of the methods applying electron bombardment, which cannot be used for bulk insulators.

In the present paper the electronic structure of amorphous silicon dioxide layer on single crystalline Si was studied by using X-ray photoemission and, for the first time, X-ray inverse photoemission method.

2. Experimental

The $\alpha\text{-SiO}_2$ sample was prepared in the Institute of Electronic Materials Technology in Warsaw. The SiO_2 film was deposited on a Si(111) substrate by rf sputtering. The film thickness was about 3000 Å. The film obtained was not stoichiometric; the composition $\text{SiO}_{1.85}$ of the layer was established by using an optical reflectivity method. For such a concentration the lattice is essentially that of $\alpha\text{-SiO}_2$ with Si as an impurity [3].

The film thickness of 3000 Å was chosen as being relatively high compared with the mean free path 50 Å for 5400 eV electrons and, on the other hand, relatively low in comparison with the electron penetration depth 0.5 μm for 5400 eV electrons in SiO_2 .

The X-ray inverse photoemission studies were performed by using an isochromat method. In this method an X-ray monochromator was set to pass a narrow band of radiation at a fixed frequency. The intensity of bremsstrahlung radiation was recorded as a function of voltage on the X-ray tube. A Johann monochromator, with Si(220) bent crystal was set at the photon energy 5415 eV. The total spectrometer resolution was 1.8 eV. The vacuum in the X-ray tube was of the order of 10^{-6} torr. The BIS spectra of $\alpha\text{-SiO}_2$ layer, as well as for Si crystal and bulk Pd, were measured with 0.5 eV/1 min or 0.5 eV/2 min steps. BIS measurements were performed for a couple of $\alpha\text{-SiO}_2$ samples to optimize experimental conditions. Several isochromats of each sample were summarized.

3. Results and discussion

In Fig. 1 the measured isochromats of $\alpha\text{-SiO}_2$ film and Si crystal are shown. The BIS intensity is very low, namely 2200 counts/26 min and 600 counts/8 min at peak maximum for $\alpha\text{-SiO}_2$ and Si isochromats, respectively. Both isochromats are very different in shape. A strong and broad peak of unoccupied electron states at about 10 eV above the threshold is seen in $\alpha\text{-SiO}_2$ isochromat. The zero energy in the $\alpha\text{-SiO}_2$ isochromat was fixed at the Fermi level, which was determined from an isochromat of gold evaporated as a 120 Å layer onto the surface of silicon dioxide.

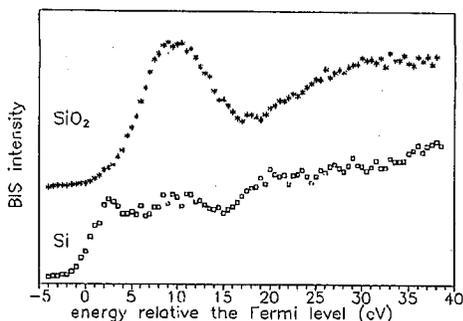


Fig. 1. X-ray bremsstrahlung isochromats of a-SiO₂ and Si-crystal samples at the photon energy 5415 eV.

Within the accuracy of 0.2 eV, the energy of Fermi level (E_F) of silicon dioxide was found to be equal to E_F of the massive Pd. Therefore we concluded that a charge shift of BIS spectrum for our a-SiO₂ sample was negligibly small.

A big problem was a decomposition of silicon dioxide during electron bombardment. Without any special thermal contact between the silicon substrate and the water-cooled sample holder, our oxide layer disappeared in twenty hours of electron bombardment at a typical bias of the X-ray tube equal to 5.4 kV and 18 mA of a beam current. A stripping of an SiO₂ film from the Si substrate during electron bombardment was clearly seen in the shape of BIS spectra, which changed continuously from a typical SiO₂ shape to a form characteristic of Si. We found that a decomposition of a-SiO₂ can be effectively reduced by a good thermal contact between the sample and holder, as well as by decreasing the current of the X-ray tube down to 11 mA. At such conditions the BIS spectrum of a-SiO₂ remained almost unchanged through 140 hours of electron bombardment.

In order to estimate a decomposition rate of SiO₂ during BIS measurements, X-ray photoelectron spectra of O 1s and Si 2p_{3/2} core levels were studied for the SiO₂ layer before and after BIS measurements. X-ray photoelectron measurements were made in a VG ESCALAB 210 spectrometer. We used the Mg K_{α} excitation ($h\nu = 1253.6$ eV) from an X-ray tube (15 kV, 20 mA). The surface was cleaned by argon ion bombardment of 2.5 kV and 5 μ A during 5 min. We assumed that such sputtering does not introduce any decomposition of the oxide. The C 1s, O 1s and Si 2p core levels and valence band photoelectron spectra were recorded before and after ion sputtering.

The XPS O 1s and Si 2p_{3/2} core level spectra together with fitting curves are shown in Figs. 2 and 3. We found that all XPS spectra are shifted to higher binding energies by about 2 eV due to charge effect. The true binding energies were determined, with respect to the Fermi level, assuming the O 1s binding energy in SiO₂ to be equal to 533.3 eV [3]. The results are presented in Table. We estimated the overall accuracy of the XPS measurement to be 0.2 eV.

The Si 2p_{3/2} core-level spectrum of the investigated a-SiO₂ layer (Fig. 3a) was fitted to a single Gaussian-like peak with its FWHM being only 0.1 eV higher than the width of the O 1s peak. Therefore, we can conclude that a lattice of a-SiO₂

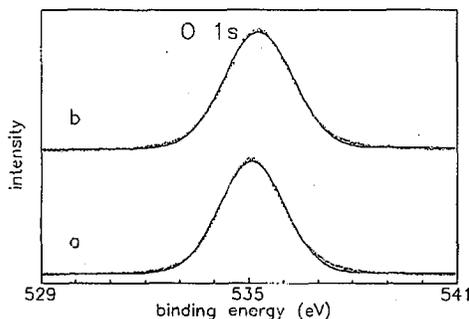


Fig. 2. O $1s$ photoelectron spectra for a-SiO₂ layer: (a) before BIS measurements, (b) after BIS measurements. Experimental data are shown by points and fitting Gaussians are drawn by solid lines. A charge shift of about 2 eV to higher binding energy is visible.

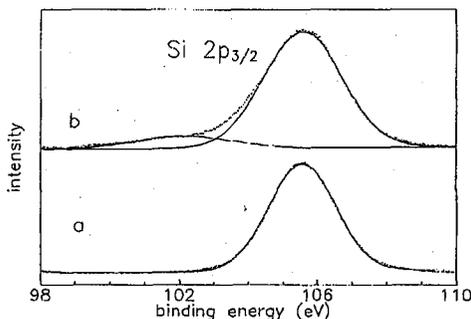


Fig. 3. Si $2p_{3/2}$ photoelectron spectra for a-SiO₂ layer: (a) before BIS measurements, (b) after BIS measurements. Experimental data are shown by points and fitting Gaussians are drawn by solid lines. A charge shift of about 2 eV to higher binding energy is visible.

TABLE
XPS binding energies and FWHM's of core levels (in eV) for a-SiO₂ before and after long-time BIS measurements.

	O $1s$		Si $2p_{3/2}$			
	E_B	FWHM	E_B	FWHM	E_B	FWHM
Before BIS	533.3	2.2	103.8	2.3	—	—
After BIS	533.3	2.3	103.6	2.7	100.1	3.2

studied consists of tetrahedral bonding SiO₄ units with small amount of Si-SiO₃ tetrahedral units of a suboxide. Another Si $2p_{3/2}$ spectrum was observed after long-time electron bombardment (Fig. 3b), where additional small peak (about 10% of the total integral intensity) appears at a binding energy shifted 3.5 eV

downward the main peak (compare with Fig. 3a). This broad peak can be attributed to Si-Si₄ and Si-Si₃O tetrahedra [3]. We also obtained that the main peak Si 2p_{3/2} became wider by 0.4 eV and shifted to the lower binding energy by 0.2 eV (see Table) after electron bombardment. It can be ascribed to the increasing concentration of Si-SiO₃ tetrahedra in comparison with that of SiO₄ tetrahedra.

The XPS valence band spectrum of a-SiO₂ film was measured in a binding energy range from 40 to -5 eV, with 0.05 eV/0.2 s steps, through 150 scans. The maximal intensity in this region was 120 000 counts/30 s. The XPS valence band spectrum was corrected for Mg K_{α3,4} satellites and charge shift. The corrected photoemission valence band spectrum and inverse photoemission spectrum were put together in a common energy scale in Fig. 4. A photoelectron spectrum of a-SiO₂

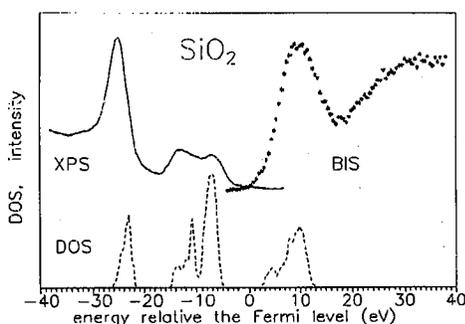


Fig. 4. Valence band photoemission and X-ray inverse photoemission spectra of a-SiO₂ on a common energy scale as discussed in the text. Below the theoretical DOS of a-SiO₂ [20] is shown.

displays three clearly discernible features. The peak at binding energy of about -25 eV refers to the O 2s orbital. The O 2p-Si 3p, 3s bonding orbital appears at binding energy -13 eV, and the O 2p nonbonding orbital is found at binding energy -5 eV at the valence band top. The inverse photoemission spectrum displays two broad maxima at 10 and 30 eV. In molecular approach the first one is an O 2p-Si 3s, 3p antibonding orbital, while the second one is discussed below as σ -resonance. In Fig. 4 the theoretical total density of electron states for a-SiO₂ calculated by Ching [20] is also shown. An energy gap of about 9 eV between the photoemission and inverse photoemission peaks is visible. The distance 17 eV between maxima of those peaks for a-SiO₂ layer studied here is in good agreement with theoretical model for amorphous SiO₂ [20].

The XPS emission from O 2s level is strong in comparison with those peaks from higher-lying O 2p orbitals. Also the XPS intensity of O 2p nonbonding orbital seems relatively smaller than it could result from the calculations of DOS (Fig. 4). These amplitude modifications are due to the energy dependent transition probabilities [21], which weight the components of the photoemission spectrum. Thus, the observed spectrum does not reproduce the true density of states, but it only shows a weighted density of states.

The energy dependence of BIS transition probability closely resembles that of

photoemission: the l -partial transition probability decreases with increasing photon energy and the rate of decrease becomes larger with the l number [22]. We estimated that relative probabilities of X-ray bremsstrahlung transitions to final electron states O $3s$, O $3p$, Si $3s$, Si $3p$ and Si $3d$ for photon energy 5415 eV are equal to 1/6, 1/100, 1, 1/9 and 1/1000, respectively. For the estimation we used data calculated in Coulombic approximation by Hahn et al. [23] and a scaling procedure proposed by these authors. Therefore, the probability of BIS transitions to Si s -like states located at the bottom of conduction band of SiO₂ is higher than this of transitions to Si p -like states placed few eV above the bottom. On the other hand, the O s -like and O p -like contributions to BIS intensity are small in the range studied due to low local density of unoccupied electron states around oxygen ions in comparison with silicon ions, according to theoretical models, which predict that the occupied electron states are localized at oxygen ions, while the unoccupied electron states are at the silicon ions [20, 24]. Therefore, in the first approximation the BIS spectrum of SiO₂ should reflect the total density of unoccupied states with an enhancement of Si s -type states contributions.

A second BIS maximum appears at the energy of 30 eV above the Fermi level, i.e. about 20 eV above the first BIS maximum. This energy distance agrees well with the energy of plasmon loss in SiO₂ ($E_p = 21$ eV) evaluated from the XPS spectra. We suggest that the second maximum originates rather from the electronic structure than from a plasmon loss, because a similar peak was observed in the absorption spectra at about 20 eV above the "white" line [9, 10, 12, 14] and was interpreted as a σ -resonant scattering of photoelectrons by adjacent atoms [14, 25].

In conclusion, a complementary study of both occupied and unoccupied electron states in amorphous SiO₂ by using X-ray photoemission (XPS) and, for the first time, by X-ray inverse photoemission (BIS) has been performed. We have aligned the photoemission valence band spectrum with BIS spectrum in a common energy scale. Our experimental results are consistent with the model calculations of density of states for a-SiO₂ by Ching [20]. We have performed an analysis of the Si $2p$ and O $1s$ core-level spectra for samples before and after BIS measurements, which showed a partial decomposition of SiO₂ due to a long-time electron bombardment during BIS experiment.

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