RARE-EARTH EXCITATION MECHANISM IN WIDE BAND GAP II–VI COMPOUNDS

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The excitation mechanism of rare-earth emission in Eu and Ce doped CaS and SrS is studied. It is proposed that the Eu and probably also Ce emission is induced by the photoionization transition of the rare-earth ion, which is followed by the carrier trapping via the excited state of the ion. At increased temperatures the efficiency of excitation is reduced. We explain this effect by the carrier emission from the excited core state of the rare-earth ion to the continuum of the conduction (valence) band states. It is also suggested that the charge transfer state of the rare-earth ion may act as the intermediate state in the carrier trapping.
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1. Introduction

Rare-earth (RE) activated wide band gap CaS and SrS belong to the most perspective candidates for producing efficient multicolor thin film electroluminescence (EL) displays [1]. Even though there are still problems with their chemical instability and hygroscopic nature, CaS and SrS doped with Eu and Ce can be applied for red (CaS:Eu) and blue (SrS:Ce) EL displays [2]. The main hindrance for obtaining the high EL efficiency from such displays was low excitation efficiency. This limitation can be avoided if the impact excitation of the RE ion is replaced by the impact ionization process [3]. This idea was verified in our recent studies of Yb doped ZnS and Eu doped ZnS and CaS [4–8]. It was shown that the efficient Eu$^{2+}$ emission can be induced by the Eu ionization. This was explained by trapping of the ionized carrier via the $4f^55d^1$ excited state of the ion ($\text{Eu}^{3+} + e \rightarrow \text{Eu}^{2+,*}$)
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[9, 10]. Different situation was found for the Eu ion in the ZnS [8]. There, the carrier trapping proceeds via the extended excited state of the RE ion (charge transfer state [11]), which is the lowest excited state of the Eu\(^{2+}\) ion in the ZnS [8]. Thus, carrier trapping by the ionized RE ion can be followed by either Eu core emission or by the so-called localized bound exciton emission [11]. The relevant importance of these two recombination processes is analyzed in this paper.

2. Results and discussion

The CaS and SrS samples studied were grown by the atomic layer epitaxy (ALE) method [12] and were doped with Eu to 0.14 w% (weight%) (CaS), 0.2 w% (SrS) and with Ce to 5.8 w% (SrS) level. Photoluminescence (PL) spectrum was excited with the Inova 200 argon laser, dispersed with the GDM 1000 double grating monochromator, and was detected with the Si EMI photomultiplier and the Stanford Research SR530 lock-in amplifier. The samples were mounted in Leybold close cycle cryogenic refrigerator.

![Fig. 1. The photoluminescence spectra of the ALE grown CaS:Eu (a), SrS:Eu (b) and SrS:Ce (c) samples measured at 10 K and 300 K. We tentatively assign two-emission bands observed in the SrS:Eu sample to the 4f\(^6\)5d\(^1\) \rightarrow 4f\(^7\) transition of the Eu ion in two different lattice sites.](image)

The measured PL spectrum, shown in Fig. 1a–c, consists of broad red (CaS, SrS:Eu) and blue (SrS:Ce) emission lines. Large width of these PL spectra is caused by inhomogeneous structure of thin films. The PL spectrum of the SrS:Eu sample shows a two-band structure, which was not observed for the samples with the larger Eu concentration [13].

Figure 2a, b shows the temperature dependence for Eu doped CaS and SrS samples. In both cases the resulting RE emission can be thermally deactivated. The observed decrease in the PL intensity with increasing temperature can be described by [9]

\[
I(T) = \frac{I(0)}{1 + A \exp(-\Delta E/kT)}.
\]
The above formula well describes the temperature dependence of the SrS:Ce and SrS:Eu (higher energy PL) emissions with $\Delta E = 6$ meV (SrS:Ce) and $\Delta E = 52$ meV (SrS:Eu). However, the CaS:Eu and SrS:Eu (lower energy band) data are better described by the two-exponential dependence with $\Delta E_1$ and $\Delta E_2$ equal to 9 meV and 50 meV for Eu ion in CaS, and 12 meV and 90 meV for SrS:Eu, respectively.

The excitation energy of Eu emission was selected within the broad photoionization bands. Thus, Eu ions are ionized by the incident light and the subsequent carrier trapping must occur for the Eu$^{2+}$ excitation. When carrier is trapped, it can be thermally ionized back to the continuum states, which deactivates PL. The PL changes can be described by the formula (1). However, we found that for CaS:Eu and for lower energy band in SrS:Eu the formula (2) better describes the experimental results. This suggests that in these two cases the carrier trapping is a two-step process. Carrier is first trapped to some intermediate state from which it can relax to the $4f^65d^1$ state of the Eu$^{2+}$ ion or can be thermally excited to the continuum states. The nature of intermediate state is not known. However, by analogy to our previous results for Eu and Yb doped ZnS [7, 8], we propose that it is the charge transfer state of the Eu ion. Following this idea, we tentatively ascribe the larger of the two deactivation energies to the energy distance between the $4f^65d^1$ excited state of the Eu$^{2+}$ ion and the edge of the CaS, SrS conduction band. The obtained energies depend on the Eu concentration in the sample [9, 13]. We ascribe the lower energy to the binding energy of the delocalized carrier bound at the Eu ion, i.e., to the binding energy of the charge transfer state of the ion.

The first experimental results for Ce ion in the SrS indicate that Ce$^{3+}$ emission in this material is also thermally deactivated. The similarity between $\Delta E$ and $\Delta E_1$ deactivation energies observed in the experiment suggests that the $\Delta E$
energy describes the carrier emission from the charge transfer state of the Ce ion in SrS.

References