

ULTRASONIC ATTENUATION IN ALLOYS

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Ultrasonic attenuation was evaluated in metallic alloys, $\text{Ni}_x\text{Cu}_{1-x}$ ($x = 1.00, 0.70, 0.60, 0.45$ and 0) due to phonon-phonon (p-p) interaction and thermoelastic loss in a wide temperature region along $\langle 110 \rangle$ crystallographic direction for longitudinal and shear waves. Ultrasonic Grüneisen parameters, nonlinearity constants and ultrasonic attenuation due to p-p interaction and thermoelastic loss were determined from 50 K to 500 K using the Born-Mayer and electrostatic potentials. The results were compared with available experimental results.

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1. Introduction

Ultrasonic attenuation studies have been recently made in metals [1-3], semiconductors [4, 5] and dielectric materials [6, 7] as a function of temperature, frequency and direction of propagation. From all these studies it is well established that at higher temperatures, i.e. 50 K and above, phonon-phonon interaction is the dominant cause of the ultrasonic attenuation in all types of solids. In metals, due to higher thermal conductivity, a part of total ultrasonic attenuation arises due to thermoelastic loss, too. In the present investigation we made theoretical evaluation of the ultrasonic absorption coefficient (upon frequency square) in Ni, $\text{Ni}_{70}\text{Cu}_{30}$, $\text{Ni}_{60}\text{Cu}_{40}$, $\text{Ni}_{45}\text{Cu}_{55}$ and Cu as a function of temperature along $\langle 110 \rangle$ direction for longitudinal and shear waves, and shear waves polarized along $\langle 001 \rangle$ and $\langle 1\bar{1}0 \rangle$ directions. A model potential consisting of electrostatic and Born-Mayer [8] type potentials was used to evaluate second- and third-order elastic constants (SOEC and TOEC) of these alloys and metals, using nearest-neighbour distance and hardness parameter of the substance as input data.

2. Theory

2.1. Theory of elastic constants

The potential used for evaluation of SOEC and TOEC is

$$\phi = \phi_c + \phi_r, \quad (1)$$

TABLE I

Expressions for SOEC and TOEC at absolute zero.

$$\begin{aligned}
C_{11}^0 &= -0.87378(e^2/r_0^4) + (1/\rho r_0)(1/r_0 + 1/\rho)\phi(r_0) \\
&\quad + (2/\rho r_0)(\sqrt{2}/2r_0 + 1/\rho)\phi(\sqrt{2}r_0) \\
C_{12}^0 &= C_{44}^0 = 0.347775(3/2)(e^2/r_0^4) + (1/\rho r_0)(\sqrt{2}/2r_0 + 1/\rho)\phi(\sqrt{2}r_0) \\
C_{111}^0 &= 10.2639(e^2/r_0^4) - (1/\rho)(3/r_0^2 + 3/\rho r_0 + 1/\rho^2)\phi(r_0) \\
&\quad - (1/2\rho)(3\sqrt{2}/r_0^2 + 6/\rho r_0 + 2\sqrt{2}/\rho^2)\phi(\sqrt{2}r_0) \\
C_{112}^0 &= C_{166}^0 = 1.209625(e^2/r_0^4) + (1/4\rho)(3\sqrt{2}/r_0^2 + 6/\rho r_0 + 2\sqrt{2}/\rho^2)\phi(\sqrt{2}r_0) \\
C_{123}^0 &= C_{144}^0 = C_{456}^0 = 0.678375(e^2/r_0^4)
\end{aligned}$$

TABLE II

Vibrational energy contribution to elastic constants.

$$\begin{aligned}
a_{11}^{\text{vib}} &= G_0(G_2 - G_1^2/6); \quad a_{12}^{\text{vib}} = G_0(G_{1,1} - G_1^2/6); \quad a_{44}^{\text{vib}} = G_0G_{1,1}/4 \\
a_{111}^{\text{vib}} &= G_0(G_3 - G_2G_1/2 + G_1^3/18) \\
a_{112}^{\text{vib}} &= G_0(G_{2,1} - G_{1,1}G_1/3 - G_2G_1/6 + G_1^3/18) \\
a_{123}^{\text{vib}} &= G_0(G_{1,1,1} - G_{1,1}G_1/2 + G_1^3/18) \\
a_{144}^{\text{vib}} &= G_0(G_{1,1,1} - G_{1,1}G_1/2) \\
a_{166}^{\text{vib}} &= G_0(G_{1,1,1} - G_{1,1}/2 + G_{2,1}) \\
a_{456}^{\text{vib}} &= G_0G_{1,1,1}; \quad G_{1,1,1} = 0, \quad G_0 = k/8r_0^3
\end{aligned}$$

TABLE III

Expressions for various G_n .

$$\begin{aligned}
G_3 &= 2\rho_0\phi(r_0)(30 + 30\rho_0 + 9\rho_0^2 - \rho_0^3 - \rho_0^4)/G + 2G_{2,1} \\
G_{2,1} &= \sqrt{2}\rho_0\phi(r_0\sqrt{2})(30 + 30\sqrt{2}\rho_0 + 18\rho_0^2 - 2\sqrt{2}\rho_0^3 - 4\rho_0^4)/4G \\
G_2 &= 2\rho_0\phi(r_0)(-6 - 6\rho_0 - \rho_0^2 + \rho_0^3)/G + 2G_{1,1} \\
G_{1,1} &= \sqrt{2}\rho_0\phi(r_0\sqrt{2})(-6 - 6\sqrt{2}\rho_0 - 2\rho_0^2 + 2\sqrt{2}\rho_0^3)/G \\
G_1 &= 2\rho_0\phi(r_0)(2 + 2\rho_0 - 2\rho_0^2) + 2\sqrt{2}\rho_0\phi(r_0\sqrt{2})(2 + 2\sqrt{2}\rho_0 - 2\rho_0^2)/G \\
G &= \rho_0\phi(r_0)(\rho_0 - 2) + \sqrt{2}\rho_0\phi(r_0\sqrt{2})(\rho_0\sqrt{2} - 2)
\end{aligned}$$

where $\rho_0 = r_0/\rho$, $\phi(r_0) = B \exp(-r/\rho)$, $\phi(r_0\sqrt{2}) = A \exp(-r_0\sqrt{2}/\rho)$
and k is Boltzmann constant.

where ϕ_c is electrostatic potential and ϕ_r is repulsive potential given as

$$\phi_c = z^2e^2/r \quad \text{and} \quad \phi_r = A \exp(-r/\rho). \quad (2)$$

Following Brugger's [9] definition of elastic constants, and starting from nearest-neighbour distance and hardness parameter of the substance and taking interactions effective up to second nearest-neighbour distance, SOEC and TOEC

were obtained at absolute zero as shown in Table I. According to the lattice dynamics developed by Leibfried et al. [10], lattice energy changes with temperature, hence adding vibrational energy contribution to elastic constants at absolute zero one gets C_{ij} and C_{ijk} (SOEC and TOEC) at required temperature as

$$C_{ij}(T) = C_{ij}^0 + a_{ij}T, \quad (3)$$

$$C_{ijk}(T) = C_{ijk}^0 + a_{ijk}T, \quad (4)$$

where superscript 0 has been used to denote SOEC and TOEC at absolute zero. a_{ij} and a_{ijk} are defined as

$$a_{ij} = l_1 k(\partial C_{ij}/\partial r)_{r=r_0} + a_{ij}^{\text{vib}}, \quad (5a)$$

$$a_{ijk} = l_1 k(\partial C_{ijk}/\partial r)_{r=r_0} + a_{ijk}^{\text{vib}}, \quad (5b)$$

where a_{ij}^{vib} and a_{ijk}^{vib} are various constants as given in Table II, various G_n appearing in Table II have been given in Table III. l_1 is a function of nearest-neighbour distance and hardness parameter and for face centred cubic crystals given as

$$l_1 = -\rho[(2 + 2\rho_0 - \rho_0^2)\phi(r_0) + 2\sqrt{2}(1 + 2\rho_0 - \rho)\phi(r_0\sqrt{2})]/l_2, \quad (6a)$$

where

$$l_2 = 2[(\rho_0 - 2)\phi(r_0) + 2(\rho_0 - \sqrt{2})\phi(r_0\sqrt{2})] \\ \times [(\rho_0 - 2)\phi(r_0) + 4(\rho_0 - \sqrt{2})\phi(r_0\sqrt{2})]. \quad (6b)$$

2.2. Ultrasonic attenuation due to phonon-phonon interaction

Ultrasonic attenuation due to phonon-viscosity mechanism (p-p interaction) [11] is given as

$$\alpha = \frac{\omega^2 \tau_{\text{th}} E_0 \Gamma}{3dv^3(1 + \omega^2 \tau_{\text{th}}^2)}. \quad (7)$$

The above expression in Akhiezer regime [12] ($\tau_{\text{th}}\omega \ll 1$) reduces to

$$(\alpha/f^2)_l = \frac{(2\pi)^2 \Gamma_l E_0 \tau_{\text{th}}}{3dv_l^3}, \quad (8)$$

$$(\alpha/f^2)_s = \frac{2\pi^2 \Gamma_s E_0 \tau_{\text{th}}}{3dv_s^3} \quad (9)$$

for longitudinal and shear waves respectively. Γ is nonlinearity coupling constant (subscripts l and s denote longitudinal and shear waves respectively), E_0 is energy density of the crystal, d is density, v_l being longitudinal wave velocity and v_s shear wave velocity and τ_{th} is thermal relaxation time given as

$$(1/2)\tau_l = \tau_{\text{th}} = \tau_s = 3k/C_v \langle v \rangle^2, \quad (10)$$

where k is thermal conductivity, C_v is specific heat per unit volume and $\langle v \rangle$ is the Debye average velocity obtained as

$$\langle v \rangle = [(1/3)(2/v_s^3 + 1/v_l^3)]^{-1/3}. \quad (11)$$

Γ is obtained as

$$\Gamma = 9\langle\gamma_i^j\rangle - \frac{3\langle\gamma_i^j\rangle^2 C_V T}{E_0}, \quad (12)$$

where $\langle\gamma_i^j\rangle$ and $\langle\gamma_i^j{}^2\rangle$ are average and square average Grüneisen numbers [11].

Ultrasonic attenuation due to thermoelastic loss $(\alpha/f^2)_{th}$ is obtained as

$$(\alpha/f^2)_{th} = \frac{2\pi^2 kT \langle\gamma_i^j\rangle^2}{dv_1^5}. \quad (13)$$

For shear waves $\langle\gamma_i^j\rangle$ becomes zero, hence there is no thermoelastic loss due to propagation of shear waves.

3. Evaluation

Taking nearest-neighbour distance [13] (Ni = 2.489, Ni₇₀Cu₃₀ = 2.516, Ni₆₀Cu₄₀ = 2.516, Ni₄₅Cu₅₅ = 2.526 and Cu = 2.558, all in Å) and hardness parameter [14] of the substances, SOEC and TOEC were obtained at different temperatures using the theory developed above. SOEC and TOEC thus obtained were used to obtain $\langle\gamma_i^j\rangle$ and $\langle\gamma_i^j{}^2\rangle$ along (110) direction for longitudinal wave over 39 pure modes and for shear waves polarized along (001) direction over 20 modes and for polarization along $\langle 1\bar{1}0 \rangle$ over 18 modes using Mason's Grüneisen table [11, 15]. C_V and E_0 have been obtained [16] as a function of Θ_D/T , at different temperatures (Θ_D is the Debye temperature) as shown in Table IV.

TABLE IV
Specific heat per unit volume, C_V (10^7 erg/cm³ K) and thermal energy density, E_0 (10^9 erg/cm³) at different temperatures for different NiCu alloys.

% of Cu in NiCu		Temperature [K]					
		50	100	200	300	400	500
0	C_V	0.38	1.64	3.06	3.40	3.55	3.62
	E_0	0.03	0.56	3.00	6.22	9.89	13.25
30	C_V	0.47	1.77	3.00	3.49	3.54	3.54
	E_0	0.03	0.62	3.25	6.35	9.77	13.76
40	C_V	0.49	1.82	3.01	3.36	3.46	3.52
	E_0	0.03	0.65	3.19	6.29	9.97	14.09
55	C_V	0.53	1.88	3.02	3.34	3.45	3.49
	E_0	0.06	0.69	3.26	6.47	9.83	14.10
100	C_V	0.70	2.09	3.04	3.29	3.37	3.41
	E_0	0.25	0.82	3.25	6.69	10.00	12.81

Nonlinearity constants (Γ_l and Γ_s) were obtained using Eq. (12). Thermal relaxation time was obtained taking thermal conductivity [16–18] with the help

TABLE V
Thermal relaxation time (10^{-11} s) and nonlinearity constant ratio at different temperatures for different NiCu alloys.

% of Cu in NiCu		Temperature [K]					
		50	100	200	300	400	500
0	τ_{th}	5.30	2.08	0.56	0.57	0.50	0.46
	Γ_1/Γ_s^*	1.16	1.42	1.49	1.51	1.55	1.54
	Γ_1/Γ_s^{**}	0.28	0.36	0.39	0.38	0.37	0.37
30	τ_{th}	6.24	5.11	2.406	2.07	2.09	2.07
	Γ_1/Γ_s^*	0.75	1.15	1.50	1.79	1.84	2.14
	Γ_1/Γ_s^{**}	0.31	0.44	0.57	0.70	0.73	0.86
40	τ_{th}	5.32	4.52	2.20	1.98	1.87	1.87
	Γ_1/Γ_s^*	0.90	1.22	1.55	1.77	1.94	21.14
	Γ_1/Γ_s^{**}	0.35	0.47	0.60	0.68	0.77	0.86
55	τ_{th}	5.86	3.90	1.82	1.69	1.59	1.57
	Γ_1/Γ_s^*	0.98	1.22	1.58	1.77	1.90	2.14
	Γ_1/Γ_s^{**}	0.31	0.46	0.60	0.68	0.74	0.85
100	τ_{th}	7.85	5.56	3.22	2.88	0.68	0.63
	Γ_1/Γ_s^*	1.04	1.20	1.34	1.38	1.37	1.16
	Γ_1/Γ_s^{**}	0.42	0.46	0.52	0.54	0.54	0.49

* For shear wave polarised along $\langle 001 \rangle$ direction.

** For shear wave polarised along $\langle 1\bar{1}0 \rangle$ direction.

of Eq. (10). Thermal relaxation time and nonlinearity constant ratio at different temperatures have been shown in Table V. Ultrasonic attenuation due to phonon-phonon interaction and thermoelastic loss was obtained at different temperatures using Eqs. (8), (9), and (13).

4. Results and discussion

Values of nonlinearity constant ratio Γ_1/Γ_s at different temperatures have been shown in Table V for different atomic concentrations of Cu in Ni. The ratio lies between 1 to 2 for shear wave polarized along the $\langle 001 \rangle$ direction and is less than 1 for the polarization along the $\langle 1\bar{1}0 \rangle$ direction, like other substances [5]. Thermal relaxation time τ_{th} (Table V) first increases with temperature and at higher temperatures it attains the saturation value as in the case of other metals [1]. From the magnitude of thermal relaxation time one can check the validity of the condition $\omega\tau_{th} \ll 1$ in the ultrasonic frequency range. (α/f^2) at different temperatures has been shown for longitudinal and shear waves in Figs. 1-5. From the

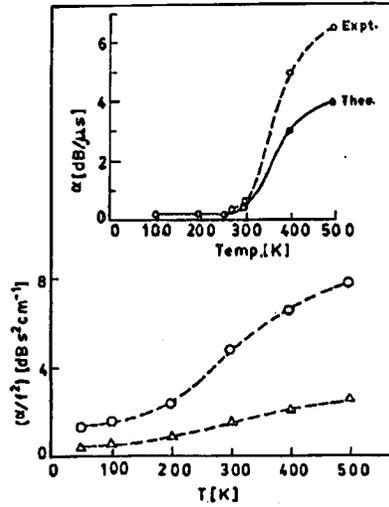


Fig. 1. Ultrasonic attenuation due to p-p interaction in Ni for shear wave (—○— in 10^{-17} dB s² cm⁻¹) polarised along $\langle 001 \rangle$ direction and for shear wave (—△— in 10^{-16} dB s² cm⁻¹) polarised along $\langle 1\bar{1}0 \rangle$ direction. Inset: experimental and theoretical values of ultrasonic attenuation for longitudinal wave in Ni.

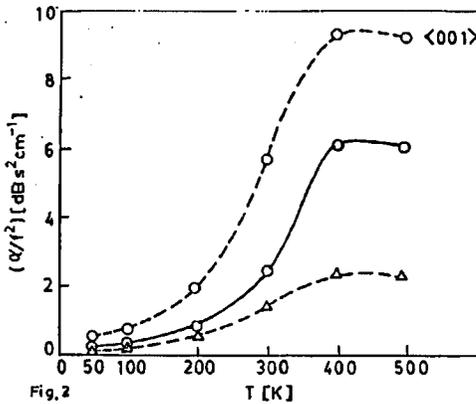


Fig. 2

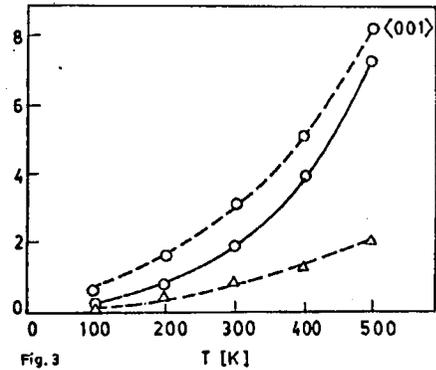


Fig. 3

Fig. 2. Ultrasonic attenuation due to p-p interaction in Ni₇₀Cu₃₀ for longitudinal wave (—○— in 10^{-16} dB s² cm⁻¹), for shear wave (—○— in 10^{-16} dB s² cm⁻¹) polarised along $\langle 001 \rangle$ direction and for shear wave (—△— in 10^{-15} dB s² cm⁻¹) polarised along $\langle 1\bar{1}0 \rangle$ direction.

Fig. 3. Ultrasonic attenuation due to p-p interaction in Ni₆₀Cu₄₀ for longitudinal wave (—○— in 10^{-16} dB s² cm⁻¹), for shear wave (—○— in 10^{-16} dB s² cm⁻¹) polarised along $\langle 001 \rangle$ direction and for shear wave (—△— in 10^{-16} dB s² cm⁻¹) polarised along $\langle 1\bar{1}0 \rangle$ direction.

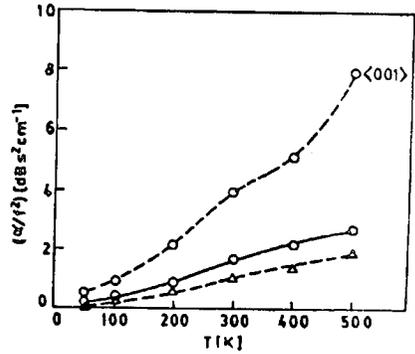
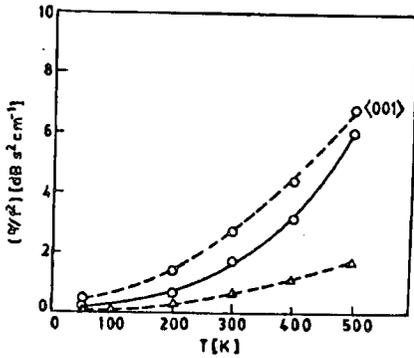


Fig. 4. Ultrasonic attenuation due to p-p interaction in $Ni_{45}Cu_{55}$ for longitudinal wave (—○— in 10^{-16} dB s^2 cm^{-1}), for shear wave (---○--- in 10^{-16} dB s^2 cm^{-1}) polarised along $\langle 001 \rangle$ direction and for shear wave (---Δ--- in 10^{-15} dB s^2 cm^{-1}) polarised along $\langle 110 \rangle$ direction.

Fig. 5. Ultrasonic attenuation due to p-p interaction in Cu for longitudinal wave (—○— in 10^{-16} dB s^2 cm^{-1}), for shear wave (---○--- in 10^{-16} dB s^2 cm^{-1}) polarised along $\langle 001 \rangle$ direction and for shear wave (---Δ--- in 10^{-15} dB s^2 cm^{-1}) polarised along $\langle 110 \rangle$ direction.

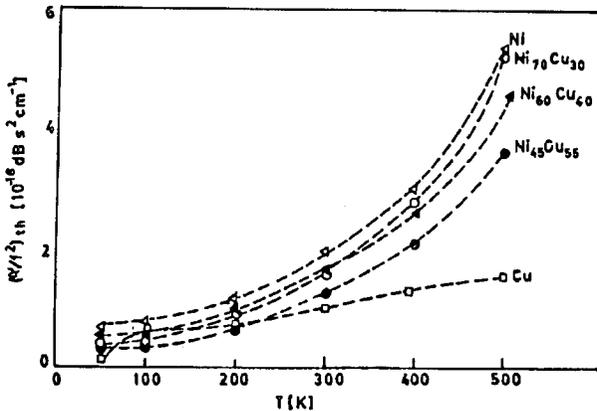


Fig. 6. Ultrasonic attenuation due to thermoelastic loss.

figures it is evident that the ultrasonic attenuation increases with temperature due to phonon-phonon interaction for all the alloys and metals studied. Experimental results are available for Ni along $\langle 110 \rangle$ direction for longitudinal wave at different temperatures at 15 MHz. We compared our results with the observed values of Basu [19] as shown in Fig. 1. From the figure one can see that the variation of the ultrasonic attenuation is the same as observed by Basu at lower temperatures.

Though no experimental results have been available for these alloys yet, we expect our values to be correct, seeing the agreement between theoretical and experimental values for Ni. The magnitude of the attenuation in the case of these metallic alloys is of the same order as observed in metals [1]. Ultrasonic attenuation due to thermoelastic loss has been shown in Fig. 6 at different temperatures. $(\alpha/f^2)_{th}$ for alloys and metals studied increases with temperature like in metals [1].

Thus, it can be concluded that the metallic alloys behave more or less like metals as far as their ultrasonic attenuation properties are concerned.

References

- [1] S.K. Kor, Kailash, *Indian J. Pure Appl. Phys.* **24**, 179 (1986).
- [2] S.K. Kor, Kailash, K. Shanker, P. Mehrotra, *J. Phys. Soc. Jpn.* **7**, 2428 (1987).
- [3] S.K. Kor, R.K. Singh, *Acustica*, submitted for publication.
- [4] S.K. Kor, Kailash, *J. Phys. Soc. Jpn.* **55**, 2232 (1986).
- [5] M.O. Manasreh, D.O. Pederson, *Solid State Ion. Diffus. React.* **15**, 65 (1985).
- [6] S.K. Kor, R.R. Yadav, Kailash, *J. Phys. Soc. Jpn.* **55**, 287 (1986).
- [7] S.K. Kor, R.K. Singh, *Acta Phys. Pol. A* **80**, 805 (1991).
- [8] M. Born, J. Mayer, *Z. Phys.* **75**, 1 (1932).
- [9] K. Brugger, *Phys. Rev. A* **133**, 1611 (1964).
- [10] G. Leibfried, W. Ludwig, in: *Solid State Physics*, Vol. 12, Eds. F. Seitz, D. Turnbull, Academic Press, New York 1961, p. 276.
- [11] W.P. Mason, *Physical Acoustics*, Vol. 3, Academic Press, New York 1965.
- [12] A. Akhiezer, *J. Phys. (USSR)* **1**, 277 (1939).
- [13] C. Kittel, *Introduction to Solid State Physics*, Wiley, New York 1981.
- [14] M.P. Tosi, in: *Solid State Physics*, Vol. 16, Eds. F. Seitz, D. Turnbull, Academic Press, New York 1964, p. 1.
- [15] W.P. Mason, R. Rosenberg, *J. Acoust. Soc. Am.* **45**, 47 (1969).
- [16] *American Institute of Physics Handbook*, Ed. D.E. Gray, McGraw-Hill, New York 1972.
- [17] C.Y. No, M.W. Ackeman, K.Y. Wus, C. Geh, T.N. Haril, *J. Phys. Chem. Ref. Data* **7**, 957 (1978).
- [18] *CRC Handbook of Chemistry and Physics*, Ed. R.C. Weast, CRC Press, Florida 1981.
- [19] A. Basu, *J. Appl. Phys.* **29**, 961 (1974).