Magnetothermal Properties of Mesoscopic System Based on Ni$_3$Pt Nanoparticle

S. Vorobiov$^a$, B. Stropkai$^{a,b}$, M. Kožejová$^a$, O. Tkach$^{a,b}$, V. Latyshev$^a$, E. Čižmár$^a$, M. Orendáč$^a$ and V. Komanický$^{a,*}

$^a$Institute of Physics, Faculty of Sciences, P.J. Šafárik University, Park Angelinum 9, 04013 Košice, Slovakia
$^b$Faculty of Electronics and Computer Technology, Suny State University, Rymskyi-Korsakov Str. 2, 40007 Sumy, Ukraine

This article discusses the influence of the geometry of mesoscopic structures based on Ni$_3$Pt nanoparticles on the magnetothermal characteristics to maximize the magnetocaloric response. Nanostructured samples in the form of arrays of nanocylinders and nanowires were made using a combination of electron-beam lithography and electrochemical deposition methods. It was determined that the change from nanoparticles to nanostructured objects contributes to an increase in the magnetocaloric effect and change of the operating temperature range.

DOI: 10.12693/APhysPolA.137.922

PACS/topics: electron-beam lithography, electrochemical deposition, magnetocaloric effect

1. Introduction

Magnetocaloric cooling is a promising alternative to classical cooling methods due to its high environmental safety and energy efficiency. In this connection, the search for new magnetocaloric materials with high magnetocaloric response in the room temperature range, free from rare-earth elements is relevant [1]. Nanomaterials are potential candidates for usage in cooling systems when compared with bulk materials because their magnetic properties significantly depend on such effects as spin-disorder effects, the interaction between neighboring atoms, grains and particles, the crystal lattice structure, the texture, and the dimensional and surface effects [2, 3]. Changing these parameters may significantly influence the functional properties of these materials. Another equally important issue is the search for appropriate geometric schemes that would maximize magnetocaloric effect (MCE) [4]. In samples with a controlled microstructure of the sample (e.g., the formation of nanowires in an ordered matrix where distance between nanowires is controlled) it is possible to implement direct and inverse MCE, which is regulated by the magnetic field [5]. This is due to the superparamagnetic behavior of nanoobjects (e.g., nanowires, nanoparticles, nanodots) and induced magnetization axes (anisotropy). The use of commercially available nanoporous membranes such as nanoporous alumina imposes nanoobject size and their geometric configuration. The solution to this problem is to use electron-beam lithography, which allows formation of arrays of nanoobjects in predetermined geometric arrangement with high accuracy. This paper presents the results of the investigations of the magnetic properties of mesoscopic structures based on Ni$_3$Pt nanoparticles of different geometries, which were obtained by electrochemical deposition.

2. Experimental methods

The fabrication of various types of arrays (nanowires or cylindrical nanomagnets) based on Ni$_3$Pt nanoparticles was carried out using electron-beam lithography and electrochemical deposition method. Electron-beam lithography was used to form masks to resist further electrochemical deposition. Initially high-resolution positive resist for direct writing of e-beam polymethyl methacrylate (NANO 950 PMMA A, MicroChem) was spin-coated on glassy carbon substrates with the size of 0.5 × 0.5 mm$^2$. At the second stage, masks were filled with a resist by electrochemical deposition. The electrolyte solution for electrodeposition consisted of 0.003 mol/l K$_2$PtCl$_4$, 0.1 mol/l NiCl$_2$, and 0.5 mol/l NaCl [6]. The dimensions of the formed arrays were 2 × 2 mm$^2$ with a period between objects of 1 µm (distance between the centers of the objects). The prepared structures were characterized by scanning electron microscopy (SEM). The estimated number of defects in the obtained arrays that may occur due to substrate imperfections and limitations in an electrochemical deposition did not exceed 0.2%. We also determined the elemental composition after electrochemical deposition by energy dispersive X-ray spectroscopy (EDS) analysis. The results show that regardless of the geometry of the structures the deviation from the specified concentration does not exceed 1%.

The study of the phase state and crystal structure was carried out using the X-ray diffraction method (X-ray) and transmission electron microscopy (TEM), respectively. It was found that both nanoparticles and various types of arrays (nanowires, cylindrical nanomagnets) based on Ni$_3$Pt have a phase composition of fcc-Ni$_3$Pt,

*corresponding author
with an average crystallite size of $d_{\text{ave}} = 8 \text{ nm}$. A more
detailed discussion of the structural-phase analysis of
structures based on Ni$_3$Pt nanoparticles was conducted
by us in [7].

Thermomagnetic measurements of the samples in the
zero-field-cooled (ZFC) and field-cooled (FC) regimes
were performed using a commercial SQUID magnetome-
ter (Quantum Design MPMS3) carried out in the temper-
ature range 2–400 K. In order to evaluate magnetocaloric
behavior, the measurement of the isothermal magnetization
curves was conducted up to 0.5 T in the temperature
range from 75 K to 400 K. The isothermal magnetic en-
tropy change $\Delta S_M$ has been calculated using the Maxwell
relation [8].

3. Results and discussion

The use of a combination of electron-beam lithogra-
phy and electrochemical deposition gave us the oppor-
tunity to form nanostructured materials with precisely
deﬁned geometry and dimensions. Figure 1 shows the
SEM image obtained from Ni$_3$Pt nanoparticles (Fig. 1a)
and fragments of microstructures with different imple-
mented geometry (Fig. 1b and c). The average size of
Ni$_3$Pt nanoparticles was determined in the previous
work [9] and was $\approx 130 \text{ nm}$. The average size of the
nanostructured objects was determined using SEM anal-
ysis. Cylindrical nanomagnets (Fig. 1b) had height of
1.7 $\mu$m with a diameter of 0.35 $\mu$m. The nanowires
(Fig. 1c) had a length of 500 $\mu$m with a height of 1 $\mu$m
and a width of 0.35 $\mu$m. It should be noted that we
consider nanowires, not in the classical sense (3D). The
investigated “nanowires” are fixed on the substrate (2D).
This configuration is more advantageous from the prac-
tical point of view since the objects of the structure can
have controlled geometry and/or location.

The study of magnetic properties showed a signifi-
cant difference in the thermomagnetic properties of Ni$_3$Pt
nanoparticles with a random orientation on the sub-
strate and structures with a precisely defined geometry.
Figure 2 shows the temperature dependence of the
zero-field-cooled/field-cooled magnetization at 10 mT ob-
tained from arrays of cylindrical nanomagnets (Fig. 2a)
and nanowires (Fig. 2b) with two measurement geomet-
ries, the field was applied parallel and perpendicular
to the matrix substrate (to the array substrate). The
blocking temperature ($T_B$), which has been determined
from the position of maximum in the temperature depen-
dence of ZFC magnetization for Ni$_3$Pt nanoparticles is
$T_B \approx 225$ K [9]. This value is significantly lower in com-
parison to the mesoscopic structures. Figure 2 shows
$T_B \approx 300$ K for an array of cylindrical nanomagnets,
while $T_B \approx 290$ K for an array of nanowires. A similar
trend in the shift of blocking temperature was observed
by other authors. For example, in [10], for Fe$_3$O$_4$ based
nanoparticles, it was shown that agglomerated nanopar-
ticles act as a cluster and have a higher $T_B$ compared to
“free” or individual nanoparticles. The temperature de-
pendence of the magnetization obtained from the arrays
of cylindrical nanomagnets remained almost unchanged
(Fig. 2a) when the measurement geometry changed from
parallel to perpendicular. A completely different situ-
ation is observed for nanowire arrays (Fig. 2b). ZFC-
FC curves have a significant difference, which can be ex-
plained by the presence of magnetic anisotropy, which
occurs due to the large length-to-diameter ratio of the
nanowires [4].

Temperature dependence of the entropy change ob-
tained for different types of structures with a field change
of 0.5 T is shown in Fig. 3. The entropy change was stud-
ied for the parallel orientation. The results of the studies
showed that the change in the entropy of the systems is
very sensitive to a change in geometry. For example, in
magnetic field of 0.5 T, for nanowires the maximum value
of $-\Delta S_M = 0.05$ J/(kg K) is observed at $T_\text{max} = 350$ K,
for cylindrical nanomagnets $-\Delta S_M = 0.13$ J/(kg K) at
360 K, and for nanoparticles Ni$_3$Pt the maximum value
of $-\Delta S_M = 0.03$ J/(kg K) is observed at 270 K.

![SEM images](image_url)
Fig. 2. (a) ZFC–FC curves obtained from arrays of cylinder nanomagnets, and (b) nanowires at a measuring field and cooling field of 10 mT.

Fig. 3. Temperature dependence of the isothermal entropy change at magnetic field 0.5 T for different geometry of the structure.

The position of the maximum in the $-\Delta S_M(T)$ dependence becomes more noticeable for the arrays of nanostructures based on Ni$_3$Pt nanoparticles, as compared to nanoparticles, and the magnitude of the effect increases. For arrays of nanostructures based on nanoparticles, the maxima in the $-\Delta S_M(T)$ dependence correspond to the value of the bulk sample, although it should be noted that the maximum has a more specific shape for arrays of nanocylinders than for arrays of nanowires, which can be explained by different height of nanostructures. As was shown in our previous work [7], for arrays of nanomagnets based on Ni$_3$Pt with the same parameters (height, chemical composition) a narrower maximum was observed for higher structures. At the same time, for nanoparticles, the maximum in temperature dependence of the isothermal entropy change has a diffused form, which is associated with the process of non-uniform blocking. This process can be explained by the wide size distribution of nanoparticles from 20 to 250 nm, which was shown in our previous work [10]. Magnetic moments of particles are blocked at certain temperatures. According to the Stoner–Wohlfarth model, the blocking temperature is defined as [11]:

$$T_B = \frac{KV}{k_B \ln(\tau_m/\tau_0)},$$

where $k_B$ is the Boltzmann constant, $V$ is the particle volume, $K$ is the anisotropy constant, $\tau_m$ is the measuring time, and $\tau_0$ is a characteristic constant of particles that it is related to gyromagnetic precession. Thus, the blocking temperature is directly proportional to the particle volume. Therefore, for nanoparticles with a certain dispersion, we see a superposition of the blocking temperature compared with nanostructured materials.

The maxima in the $-\Delta S_M(T)$ dependence correspond to the direct MCE associated with the magnetic transition from the ferromagnetic to the paramagnetic state. Although it should be noted that an inverse magnetocaloric effect is observed for Ni$_3$Pt nanoparticles at temperatures below 100 K. The presence of a direct and inverse magnetocaloric effect in systems of superparamagnetic nanoparticles was predicted by Monte Carlo simulations [12]. The calculations predict that the magnitude of inverse magnetocaloric effect depends both on the size and composition of the particles [12]. Since our sample contains certain number of particles with sizes less than 25 nm, which is very close to the critical size of Ni nanoparticles (24 nm) [13]. This might be why inverse magnetocaloric effect has been observed for this particular sample.

4. Conclusions

The arrays of mesoscopic objects with different geometries were formed on the basis of Ni$_3$Pt nanoparticles by the method of electron-beam lithography followed by electrochemical deposition. The study of the thermomagnetic characteristics showed that the change from nanoparticles to nanostructured objects shifts blocking temperature from 270 K for nanoparticles to $\approx 365$ K for nanowires and nanocylinders. The magnitude of the magnetocaloric effect also increases significantly, and the peaks in the $-\Delta S_M(T)$ dependences become more pronounced. The maximum amplitude of the MCE was observed for arrays of nanocylinders at 0.5 T, $(-\Delta S_M) = 0.13 \ J/(kg \ K)$, and 360 K. During the study of the temperature dependence of the entropy change in Ni$_3$Pt nanoparticles, the presence of an inverse magnetocaloric effect was observed.
Acknowledgments

This work has been supported by grant VEGA No. 1/0204/18, APVV-17-0059, APVV-14-0073 project PROMATECH, contract No. ITMS26220220186, and project NEMMA, contract No. ITMS: 313011T544. O. Tkach acknowledges the financial support provided by the National Scholarship Programme of the Slovak Republic (SAIA).

References