Shape Change of Micelles Dragged with Constant Velocity
as Addressed in Terms of Biolubrication Application

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Molecular dynamics simulations have been performed for a model of coarse-grained lipids in water. We try to
establish structural properties of micelles when they are dragged past each other. The lipid–lipid and lipid–water
interactions are taken from a coarse-grained model first introduced by Marrink et al. This work focuses on shape
changes upon mutual interaction of micelles when dragged with constant velocity. We find that micelles can change
their shapes from spherical to cylindrical regardless of the number of constituting lipids. We address this result
toward understanding facilitated vs. obstructed lubrication of articular cartilage by synovial fluid at nanoscale
level.

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1. Introduction and simulation method

1.1. Preliminaries of micelles-assisted lubrication

The mechanism of articular cartilage (AC) lubrication has been studied for almost a century, but still the mecha-
nism has remained uncovered. The problem of AC lubrication is connected to the complexity of the friction
phenomena in biological systems and its multiscale effects inherent in the phenomena. AC includes two oppo-
site rubbing surfaces and the synovial fluid (SF) between them. SF is composed mainly of water and amphiphilic
components such as hyaluronic acid, surface active phospholipids (SAPL), etc. According to Gadomski [1] and
refs. therein, the SAPLs in their micellar states may play a crucial role in AC lubrication [2]. Moreover, Gadomski
et al. [3] suggest that the role of micelles is to lower friction by changing a sliding friction into its rolling counter-
part represented — as reminiscent of bearings — by their spherical (or, rounded) shapes. The frictional behavior
between soft colloids is of fundamental interest, e.g., due to the potential applications of spherical star polymers or
self-assembled spherical structures in drug delivery, viscosity modifications, etc. [4]. This work presents molecu-
lar dynamics simulation studies on micellar interactions in aqueous environment, where two micelles are dragged
past each other with constant velocity. We employ dodecylphosphocholine (DPC) lipids for this simulation as
they have natural tendency to self-assemble into micelles.

1.2. Model

The simulations have been performed by adopting a model by Marrink et al. [5], i.e., by a coarse-grained (CG)
model of water–lipid interactions. The simulation box (Fig. 1) is composed of water and DPC lipids building up

![Fig. 1. Snapshot of the simulation box. Two identical micelles have been colored for better visualization (silver particles represent hydrophilic heads of lipids, whereas green and pink — hydrophobic tails). Water molecules as well as two dummy atoms (see text) are not shown in the picture. The distance between the micelles’ centers of mass in x- and y-direction are denoted as Δx and Rd, respectively. \( v_0 \) stands for drift velocity of a micelle.](image-url)
\[ V(r) = 4 \epsilon \left[ \left( \frac{\sigma}{r} \right)^{12} - \left( \frac{\sigma}{r} \right)^{6} \right], \]

as well as screened Coulomb potential are used to describe pairwise interactions between atoms/ions as described in [5]. All simulations were performed using LAMMPS [6] package. SMD (steered molecular dynamics) is a method in which a centre of mass of each micelle is attached to a dummy atom by a spring which is pulled with constant velocity. The velocity is embodied in a Peclet number defined as

\[ Pe = \frac{\tau_r}{\tau_d} = \frac{\tau_r v_d}{R_g} , \]

which for the presented simulations is equal to 0.25 (similarly to other soft colloids [4]). \( \tau_r \) stands for relaxation time, \( v_d \) is a drift velocity of a micelle and \( R_g \) is a radius of gyration, conforming to its standard definition. Micelles are pulled toward each other along \( x \)-direction and are separated in \( y \)-direction by \( \frac{R_d}{R_g} = 0.63 \). \( R_d \) is the distance between centers of mass in \( y \)-direction. The interactions between the coronas of both micelles go by employing L-J and Coulomb potentials.

2. Results

The simulations were performed for three different micelles’ aggregation numbers, namely 40, 60, and 75. Figure 2 shows shape changes of micelles in \( y \)- and \( z \)-directions in terms of two ratios \( \frac{R_y}{R_x}(\Delta x) \) and \( \frac{R_z}{R_x}(\Delta x) \), where \( R_x, R_y \) and \( R_z \) are the corresponding axial-coordinate projections of the radius of gyration \( R_g = \sqrt{R_x^2 + R_y^2 + R_z^2} \). They reveal the expectation that the smaller the aggregation number is the more mechanically stable the micelles can be. The micelles are changing their shapes as depicted in Fig. 2 from spherical to quasi-cylindrical or even pancake-like. The effect as such is not explicitly discussed in the present work. It, however, resembles fairly some nonequilibrium morphologies [7, 8], herein recoverable in terms of the aggregation number vs. drift velocity taken for different temperatures.

3. Conclusions

Our simulation results suggest that the morphology of two tightly interacting micelles depends sensitively on their mass center distance, entailed with the aggregation numbers. Implicitly, such complex dynamic behavior resembles qualitatively the nonequilibrium morphologies occurring in complex diffusive processes [7, 8]. The study may also help rationalize facilitated lubrication conditions in articulating micelles-involving systems [1, 3]. At first sight, the rationalization would rely on observing that the pancake-like structure while in motion will be more prone to deceleration as the other emergent DPC-based structures, which will make the lubrication conditions less operative [1].

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References