

Magnetization Reversal Processes in Nanocrystalline (Pr, Dy)–(Fe, Co)–B Bulk Alloys

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The aim of this paper was to study the phase constitution, magnetic properties and magnetization reversal processes in the rapidly solidified bulk (Pr,Dy)–(Fe,Co)–B alloys doped with Zr, Ti, Mn and Ni. The 3 mm outer diameter tubes samples of the $\text{Pr}_8\text{Dy}_1\text{Fe}_{60}\text{Co}_7\text{Ni}_{(6-x)}\text{Mn}_x\text{B}_{14}\text{Zr}_1\text{Ti}_3$ (where $x = 0, 3, 6$) alloys were produced by suction-casting technique. The admixture of Zr was introduced in order to improve their glass forming abilities while Dy was substituted to enhance the magnetocrystalline anisotropy of hard magnetic phase. The effect of Ni and Mn addition on the phase constitution and magnetic properties was studied in the presented work.

DOI: [10.12693/APhysPolA.127.579](https://doi.org/10.12693/APhysPolA.127.579)

PACS: 75.20.En, 75.75.-c, 75.75.Cd

1. Introduction

From the technological point of view, the amorphous and nanocrystalline Fe-based alloys are very important materials, due to their good magnetic properties and low cost [1–3]. In order to extend their commercial use, bulk components of amorphous or nanocrystalline microstructure are produced by several techniques. Relatively new procedures of obtaining bulk magnetic materials are the rapid quenching techniques — suction or injection casting methods [4, 5]. Interesting results concerning processing bulk metallic glasses were obtained for Fe–Co–Zr–W–B alloys in the form of tubes with outer diameter 3 mm and rods with maximum diameter up to 2 mm [6–8]. It was revealed that depending on the sample diameter and shape, various cooling rates can be obtained in the suction casting process [9]. This technique was also used for manufacturing nanocrystalline RE–Fe–B-type (RE = Nd, Pr, Dy) permanent magnets, where short processing time and better corrosion resistance than for magnets prepared by conventional methods seems to be advantageous. This type of products can find application as miniature ring-type magnets or stepper motors dedicated for example to electronic watches [10]. The effect of Ni and Mn addition on phase constitution and magnetic properties (Pr,Dy)–(Fe,Co)–B–Zr–Ti alloys in form of ribbons was described in our previous studies in [11]. In present paper the phase constitution and magnetic properties of $\text{Pr}_8\text{Dy}_1\text{Fe}_{60}\text{Co}_7\text{Ni}_{(6-x)}\text{Mn}_x\text{B}_{14}\text{Zr}_1\text{Ti}_3$ (where $x = 0, 3, 6$) in the form of 3 mm outer diameter tubes alloys were discussed. Furthermore, for samples containing Mn, magnetization reversal processes were investigated.

2. Samples preparations and experimental methods

Alloy ingots with nominal compositions of the $\text{Pr}_8\text{Dy}_1\text{Fe}_{60}\text{Co}_7\text{Ni}_{(6-x)}\text{Mn}_x\text{B}_{14}\text{Zr}_1\text{Ti}_3$ (where $x = 0, 3, 6$)

were prepared by arc-melting under an argon atmosphere using high purity constituent elements with pre-alloyed Fe–B of known composition. Then the samples were homogenized by systematic re-melting. Subsequently, samples in the form of 3 mm outer diameter thin walls tubes were prepared by the suction-casting technique. The phase structure of the as-cast samples was investigated by X-ray diffractometry (XRD) with Cu K_α radiation ($\lambda = 1.54$ nm). The room temperature major and minor hysteresis loops, as well as sets of recoil curves were performed using a vibrating sample magnetometer (LakeShore VSM) in the external magnetic field up to 1600 kA/m, on specimens prepared from tubes. The minor hysteresis loops were used to obtain the field dependences of remanence J_r and coercivity JH_c . The series of recoil curves were obtained for the initially saturated samples and for the thermally demagnetized specimens.

3. Results

X-ray diffraction patterns collected for $\text{Pr}_8\text{Dy}_1\text{Fe}_{60}\text{Co}_7\text{Ni}_{(6-x)}\text{Mn}_x\text{B}_{14}\text{Zr}_1\text{Ti}_3$ (where $x = 0, 3, 6$) alloys are presented in Fig. 1. All investigated samples are crystalline in the as-cast state. For samples containing 6 at.% of Ni coexistence of two phases: dominant soft magnetic α -(Fe,Ni) phase and a small amount of hard magnetic $\text{Pr}_2(\text{Fe,Co})_{14}\text{B}$ phase, were shown (Fig. 1a). As it was shown (Fig. 1b and c), the increase of Mn admixture in alloy composition leads to the increase of volume fraction of the hard magnetic $\text{Pr}_2(\text{Fe,Co})_{14}\text{B}$ phase thus leading to the improvement of the magnetic properties. Furthermore for the Mn containing samples no evidences of presence of the soft magnetic phase was shown. The peaks corresponding to recognized phases are marked in Fig. 1.

It is supposed that Mn and Ni atoms together with Zr and Ti additions were located in the disordered soft magnetic phase and/or precipitated at the boundaries crystalline phases. The Rietveld analysis carried out for sample with 6 at.% of Ni, showed larger lattice constant a of soft magnetic phase α -(Fe,Ni) than that for a pure α -Fe, which suggests that the Ni atoms occupies Fe positions in the unit cell of bcc α -Fe phase. It was also shown

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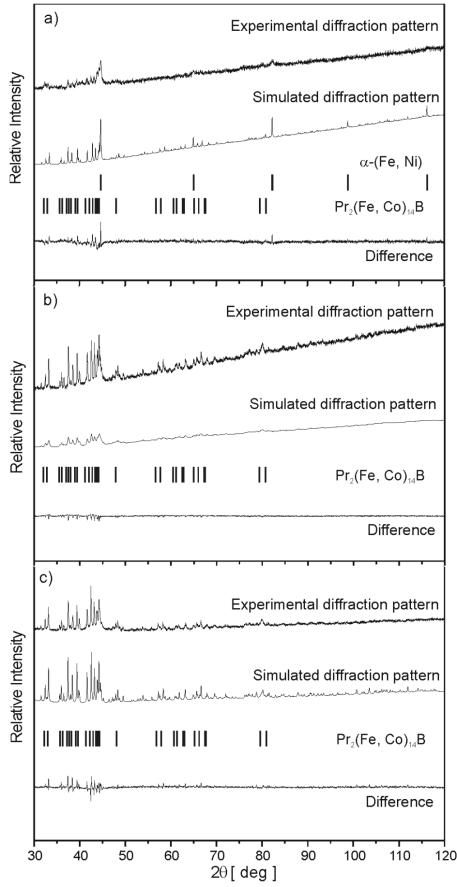


Fig. 1. The experimental X-ray diffraction data together with the Rietveld analysis and differences between experiment and calculations for the as-cast $\text{Pr}_8\text{Dy}_1\text{Fe}_{60}\text{Co}_7\text{Ni}_6\text{B}_{14}\text{Zr}_1\text{Ti}_3$ (a), $\text{Pr}_8\text{Dy}_1\text{Fe}_{60}\text{Co}_7\text{Ni}_3\text{Mn}_3\text{B}_{14}\text{Zr}_1\text{Ti}_3$ (b) and $\text{Pr}_8\text{Dy}_1\text{Fe}_{60}\text{Co}_7\text{Mn}_6\text{B}_{14}\text{Zr}_1\text{Ti}_3$ (c) alloys in the form of thin walls 3 mm outer diameter tubes.

that with the increase of Mn contents, a slight decrease of lattice constants of hard magnetic phase occurs. Such phenomenon can be related to replacement of Fe atoms by Mn in the unit cell of the in hard magnetic phase. Results of the Rietveld analysis were collected in Table I.

TABLE I

Lattice constants of recognized magnetic phases for all investigated alloys.

Sample	α -(Fe,Ni)	$\text{Pr}_2(\text{Fe,Co})_{14}\text{B}$	
	a [nm]	a [nm]	c [nm]
$\text{Pr}_8\text{Dy}_1\text{Fe}_{60}\text{Co}_7\text{Ni}_6\text{B}_{14}\text{Zr}_1\text{Ti}_3$	0.287	0.8801	1.2251
$\text{Pr}_8\text{Dy}_1\text{Fe}_{60}\text{Co}_7\text{Ni}_3\text{Mn}_3\text{B}_{14}\text{Zr}_1\text{Ti}_3$	-	0.8779	1.2179
$\text{Pr}_8\text{Dy}_1\text{Fe}_{60}\text{Co}_7\text{Mn}_6\text{B}_{14}\text{Zr}_1\text{Ti}_3$	-	0.8776	1.2167

Major hysteresis loops and initial magnetization curves measured in external magnetic field up to 1600 kA/m were shown in Fig. 2, respectively. Initial magnetization curves were used to determine saturation polarization J_s for all investigated samples. The magnetic parameters of the investigated specimens were collected in Table II.

It was shown that the highest saturation polarization had the sample doped with 6 at.% of Ni. However, the

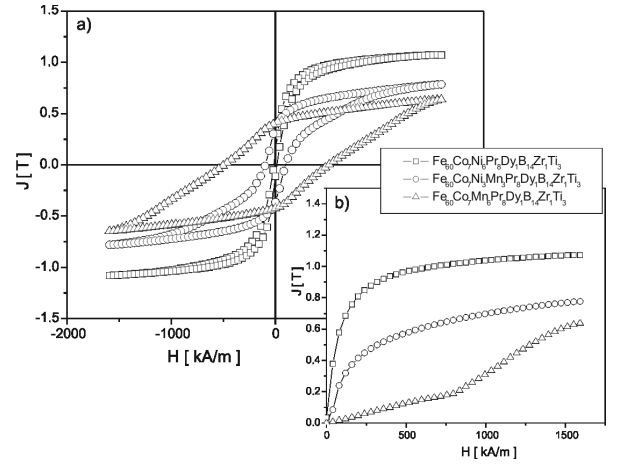


Fig. 2. The hysteresis loops (a) and initial magnetization curves (b) measured for $\text{Pr}_8\text{Dy}_1\text{Fe}_{60}\text{Co}_7\text{Ni}_{(6-x)}\text{Mn}_x\text{B}_{14}\text{Zr}_1\text{Ti}_3$ (where $x = 0, 3, 6$) alloys tubes.

TABLE II

Magnetic properties of $\text{Pr}_8\text{Dy}_1\text{Fe}_{60}\text{Co}_7\text{Ni}_{(6-x)}\text{Mn}_x\text{B}_{14}\text{Zr}_1\text{Ti}_3$ (where $x = 0, 3, 6$) alloys in the form of 3 mm outer diameter tubes.

Samples tube 3 mm as-cast	JH_c	J_r	J_s	J_r/J_s	$(BH)_{\max}$
	$[\frac{\text{kA}}{\text{m}}]$	[T]	[T]	[-]	$[\frac{\text{kJ}}{\text{m}^3}]$
$\text{Pr}_8\text{Dy}_1\text{Fe}_{60}\text{Co}_7\text{Ni}_6\text{B}_{14}\text{Zr}_1\text{Ti}_3$	15	0.121	1.08	0.11	0.399
$\text{Pr}_8\text{Dy}_1\text{Fe}_{60}\text{Co}_7\text{Ni}_3\text{Mn}_3\text{B}_{14}\text{Zr}_1\text{Ti}_3$	100	0.365	0.79	0.46	5.572
$\text{Pr}_8\text{Dy}_1\text{Fe}_{60}\text{Co}_7\text{Mn}_6\text{B}_{14}\text{Zr}_1\text{Ti}_3$	502	0.433	0.67	0.65	22.5

coercivity field JH_c measured for this alloy is relatively low just like for a semi-hard ferromagnetic. Magnetic properties were improved by the change of chemical composition of the alloy, which was made by addition of Mn atoms. With an increase of Mn admixture, an increase of remanence J_r and coercivity field JH_c was observed, while decrease of saturation polarization was noticed. The best magnetic properties were measured for sample containing 6 at.% Mn. Changes in the shape of magnetic hysteresis loops reflect changes in the phase constitution of investigated samples.

Even though the maximum magnetic field is relatively low, the obtained values of JH_c and J_s are close to maximum due to low volume fraction of the hard magnetic phase in the samples volume. Furthermore, interesting changes in the shapes of recoil curves that give insight into the magnetization reversal process occur at relatively low external magnetic fields.

Studies of magnetization reversal processes were carried out on the $\text{Pr}_8\text{Dy}_1\text{Fe}_{60}\text{Co}_7\text{Mn}_6\text{B}_{14}\text{Zr}_1\text{Ti}_3$ alloy sample revealing the best hard magnetic properties, based on measurement of recoil curves. From this measurements reversible magnetization curves (the dependences of reversible parts M_{rev} on the irreversible parts M_{irr} of magnetization for various external magnetic fields) were constructed. The shapes of M_{rev} vs. M_{irr} curves can give the information about the dominant magnetization reversal process that control the magnetic properties of the sample. For magnets where the nucleation process

is dominant, reversed domains once nucleated, can expand and easily propagate through the grains. In this case M_{rev} arises from rotation of the magnetization vectors and the $M_{\text{rev}}(M_{\text{irr}})$ plots take the shapes of straight lines of negative slopes. For magnets where the dominant magnetization reversal process is the pinning of domain walls, the $M_{\text{rev}}(M_{\text{irr}})$ curves show minimum, due to the fact that the reversible magnetization arises from the domain wall bowing [12]. Selected M_{rev} vs. M_{irr} curves for the $\text{Pr}_8\text{Dy}_1\text{Fe}_{60}\text{Co}_7\text{Mn}_6\text{B}_{14}\text{Zr}_1\text{Ti}_3$ alloy 3 mm outer diameter tube are showed in Fig. 3. The reversible magnetization curves exhibit a shallow minimum for positive values of M_{irr} , which moves toward of M_{irr} with the increase of the magnetic field. A shallow minimum on those curves suggests that for low magnetic fields the pinning of domains walls occurs, while for higher magnetic fields the nucleation process prevails. This is consistent with the shape of the initial magnetization curve (Fig. 2b) measured for the $\text{Pr}_8\text{Dy}_1\text{Fe}_{60}\text{Co}_7\text{Mn}_6\text{B}_{14}\text{Zr}_1\text{Ti}_3$ alloy sample, where for the low magnetic relatively low values of magnetic polarization were measured, while for field higher than 800 kA/m a steeper increase of magnetic polarization was measured.

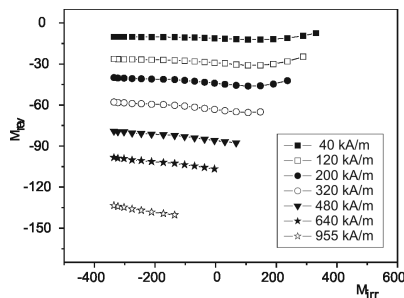


Fig. 3. Plots of reversible part of magnetization M_{rev} as a function of irreversible magnetization M_{irr} for the $\text{Pr}_8\text{Dy}_1\text{Fe}_{60}\text{Co}_7\text{Mn}_6\text{B}_{14}\text{Zr}_1\text{Ti}_3$ alloy tubes.

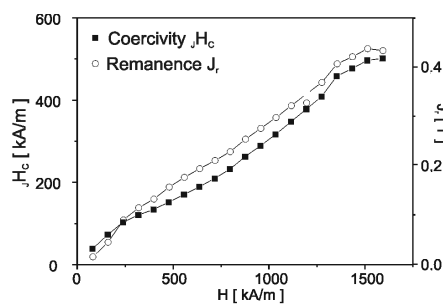


Fig. 4. Field dependences of coercivity JH_C and remanence J_r for the $\text{Pr}_8\text{Dy}_1\text{Fe}_{60}\text{Co}_7\text{Mn}_6\text{B}_{14}\text{Zr}_1\text{Ti}_3$ alloy tubes measured on initially demagnetized samples.

Additional information about the magnetization reversal process can provide the measurements of minor hysteresis loops for samples initially demagnetized. Investigation of minor hysteresis loops were carried out at room temperature using a vibrating sample magnetometer in the external magnetic field up to 1600 kA/m. From those measurements the dependences of JH_C and J_r as the function the maximum external magnetic field H_{max} were cal-

culated and presented in Fig. 4. The characteristic almost straight line shapes of this dependences confirms the results obtained from the reversible magnetization curves. For the investigated alloy, magnetization reversal process is controlled both by the reversible movement of pinned domain walls (domain walls bowing) and the nucleation of reversed domains.

4. Conclusions

It was shown that all investigated $\text{Pr}_8\text{Dy}_1\text{Fe}_{60}\text{Co}_7\text{Ni}_{(6-x)}\text{Mn}_x\text{B}_{14}\text{Zr}_1\text{Ti}_3$ (where $x = 0, 3, 6$) alloys in the form of 3 mm outer diameter tubes are crystalline in as-cast state. For the $\text{Pr}_8\text{Dy}_1\text{Fe}_{60}\text{Co}_7\text{Ni}_6\text{B}_{14}\text{Zr}_1\text{Ti}_3$ alloy the XRD analysis revealed multiphase constitution consisting of the hard magnetic $\text{Pr}_2(\text{Fe},\text{Co})_{14}\text{B}$ and soft magnetic $\alpha\text{-(Fe,Ni)}$ phases. Replacement of Ni by Mn led to growth of the hard magnetic $\text{Pr}_2(\text{Fe},\text{Co})_{14}\text{B}$ phase and improvement of magnetic parameters. For specimens containing Mn the soft magnetic phase is not observed. The Rietveld analysis of the XRD scan for the $\text{Pr}_8\text{Dy}_1\text{Fe}_{60}\text{Co}_7\text{Ni}_6\text{B}_{14}\text{Zr}_1\text{Ti}_3$ alloy has shown difference in the lattice constant of the $\alpha\text{-(Fe, Ni)}$ phase ($a = 0.287$ nm) comparing to the pure $\alpha\text{-Fe}$. In the Mn containing alloys the admixture elements are most likely located in the grain boundaries.

Studies of the magnetization reversal process for $\text{Pr}_8\text{Dy}_1\text{Fe}_{60}\text{Co}_7\text{Mn}_6\text{B}_{14}\text{Zr}_1\text{Ti}_3$ alloy tubes shown that the mechanism of magnetization reversal is controlled by both: nucleation and pinning processes. The dependences of JH_C on the maximum external magnetic field demonstrate steady rise of its value with the increase of H_{max} . Also the shapes of the M_{rev} vs. M_{irr} suggest combined magnetization reversal process.

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