Morphology, Optical and AC Electrical Properties of Copper Phthalocyanine Thin Films

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AC electrical properties of sandwich devices composed of thermally evaporated thin films of copper phthalocyanine (CuPc) with aluminum and gold electrodes (Al/CuPc/Au) are investigated over frequency (f) range of 10^2-10^5 Hz and temperature range of 293–453 K. Morphology of the samples was studied via field emission scanning electron microscope images and X-ray diffraction micrographs. The X-ray diffraction micrograph indicates the configuration of α -CuPc with the (510) plane as the preferred orientation. UV–Vis absorption spectrum was analyzed and the optical band-gap energy of CuPc thin film was determined to be 2.81 ± 0.01 eV. Capacitance increased with increasing temperature especially for $f = 10^2$ Hz. Loss factor decreased considerably with increasing frequency to a minimum value at about $f = 10^4$ Hz and increased afterwards. Capacitance is generally independent of frequency for $T \leq 413$ K; however it decreases remarkably with increasing frequency for T > 413 K. The conductivity increases quite noticeably with increasing frequency particularly for $T \leq 413$ K. The AC electrical characteristics are in good agreement with Goswami and Goswami model. According to our data, at high temperatures, the band theory is applicable in describing the conduction process, whereas hopping mechanism is dominant at low temperatures.

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1. Introduction

Metal phthalocyanines (MPcs) are organic semiconducting materials which have been widely employed as active layers in gas sensing applications over the last two decades [1–3]. They have been also used in photovoltaic devices [4] and nonlinear optics [5]. Organic semiconductor thin films exhibit a high potential for production of efficient, low cost, and flexible electronic and optoelectronic devices with the option for large area applications [6]. In this type of materials, the π -orbitals overlap of adjacent organic molecules contributes to the electrical transport [7].

Halogenated Pcs, composites of conductive polymers and Pcs, and heterojunctions of Pcs have been studied for their gas sensing properties on exposure to different gases such as NO₂, CO₂, and O₂ [8–10].

Copper phthalocyanine (CuPc) is a type of MPc whose structural [11], electrical [12] and gas sensing [2, 13–15] properties have been under intense investigation. Hassan and Gould [16] studied the effect of oxygen exposure and annealing on the conductivity of CuPc thin films and reported a discrete trap level at higher voltages, which was attributed to oxygen incorporation within the films.

In this study we intend to investigate the AC electrical characteristics of CuPc sandwich devices with Al and Au electrodes before and after exposure to oxygen on which little work has been reported so far. We study the annealing effect on the electrical properties of the devices, as well. The optical properties and morphology of samples are also analyzed using UV–Vis absorption spectroscopy, scanning electron microscopy (SEM), and X-ray diffraction (XRD) analyses.

2. Experimental

Sandwich devices were fabricated using thermal evaporation method under a vacuum of approximately 10^{-5} mbar with the evaporation rate of 0.8 nm per second. First, 70 nm of gold was evaporated on a glass substrate to act as the lower electrode. 1 μ m of CuPc was then evaporated on the gold layer. Finally 70 nm of aluminum was evaporated onto the CuPc layer to act as the upper electrode. The sandwich devices produced in this way had an area of $10 \times 10 \text{ mm}^2$.

The AC electrical measurements were then taken in the dark condition via MT4080A Motech LCR multifrequency meter over the temperature range of T = 293to 453 K and the frequency range of 10^2-10^5 Hz before and after exposure to oxygen with pressure of 0.5 mbar. We investigated the morphology and crystalline structure of CuPc thin films by field emission scanning electron microscope (FESEM) images and XRD patterns using Hitachi, Japan S-4160 and INEL, France Equinox 3000 equipment, respectively. We also studied the optical properties of CuPc thin films via Perkin Elmer, Lambda 25 UV/Vis Spectrometer.

3. Results and discussion

3.1. Structural characteristics

FESEM images of CuPc thin films with different magnifications are shown in Fig. 1a and b. The estimated

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average size of the particles is 90–100 nm and a quite uniform distribution of the roughly spherical-shaped particles is observed in Fig. 1a and b.

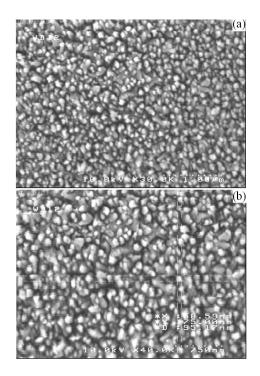


Fig. 1. FESEM images of the surface of CuPc thin film at (a) 1 μ m and (b) 750 nm magnifications indicating the mean particle diameter to be 90–100 nm.

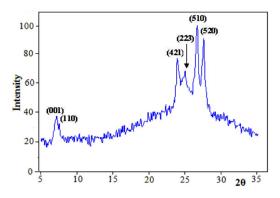


Fig. 2. X-ray diffraction micrograph of α -CuPc thin film with the thickness of 1 μ m.

The X-ray diffraction pattern of the CuPc thin films is illustrated in Fig. 2. The XRD micrograph in Fig. 2 indicates the configuration of α -form with the (510) plane as the preferred orientation. The numerical Ref. [17] and measured data are presented in Table I. The reference and measured intensity values are noticed to be different particularly for the first two peaks. Discrepancies in the relative peak intensities between conventional diffraction and X-ray diffraction patterns are largely due to the texture effect [18]. The mean crystalline size was calculated

TABLE I

X-ray diffraction data for α -CuPc sample at 25 ± 1 °C^{*a*}.

| $h \ k \ l$ | $2\theta_{\rm r}$ | $2\theta_{\rm m}$ | $I_{\rm r}$ | Im |
|-------------|-------------------|-------------------|-------------|-------|
| $0 \ 0 \ 1$ | 6.906 | 7.005 | 100 | 39.58 |
| $1 \ 1 \ 0$ | 7.246 | 7.353 | 90 | 33.75 |
| $4\ 2\ 1$ | 24.184 | 23.938 | 70 | 72.78 |
| $2\ 2\ 3$ | 25.522 | 25.025 | 70 | 64.73 |
| $5 \ 1 \ 0$ | 26.210 | 26.619 | 90 | 100 |
| $5\ 2\ 0$ | 27.616 | 27.510 | 90 | 90.94 |

^a Miller indices and diffraction angles (Ref. [17] (r) and measured (m) for α -CuPc) are given.

The intensities $I_{\rm r}$ and $I_{\rm m}$ refer to the peak height.

by Scherrer's formula [19]:

$$L = \frac{k_{\rm S}\lambda}{\beta_0 \cos\theta},\tag{1}$$

where λ is the X-ray wavelength (1.540 Å), β_0 — the full width at half maximum (FWHM) of the most intense peak in radians, θ — Bragg's angle and $k_{\rm S}$ — the Scherrer's constant = 0.9 [19]. The mean crystalline size was calculated to be about 82 nm which is in quite good agreement with the estimated value obtained from FE-SEM images.

3.2. Optical properties

The optical absorption spectrum of CuPc thin films is illustrated in Fig. 3. The UV–Vis spectrum for phthalocyanines originates from molecular orbitals within the aromatic 18π electron system and from overlapping orbitals on the central metal [20]. The conjugated double bonds within the crystal structure of the film create electron orbitals which overlap between molecular π orbitals. These electrons are able to transfer energy throughout the structure and are responsible for the absorption peaks [21].

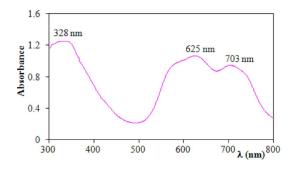


Fig. 3. UV–Vis absorption spectrum of CuPc thin film with the thickness of 1 μ m.

A doublet of peaks exists in the visible region in Fig. 3 at 625 and 703 nm known as the Q-bands. In the near UV region a single peak is observed at 328 nm (B-band) and it seems that there are two trap levels with 1.76 and 1.99 eV. The absorption coefficient α can be calculated using the relation below

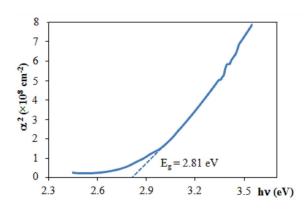


Fig. 4. α^2 as a function of photon energy $(h\nu)$ indicating a band gap of 2.81 ± 0.01 eV for CuPc thin films.

$$\alpha = 2.303 \frac{A}{t},\tag{2}$$

where A is the absorbance of the film and t is its thickness. For direct allowed transition, the absorption coefficient α is related to the photon energy $(h\nu)$ by the following relation:

$$\alpha = \alpha_0 (h\nu - E_g)^{1/2},\tag{3}$$

where $E_{\rm g}$ is the optical bandgap and α_0 is a constant [22]. In Fig. 4, α^2 is plotted as a function of photon energy $(h\nu)$ which indicates the existence of a direct optical bandgap of 2.81 ± 0.01 eV for CuPc thin films.

3.3. AC measurements 3.3.1. Frequency and temperature effects on conductivity

Figure 5 shows the dependence of conductivity (σ) on frequency (f) at different temperatures (T). The conductivity is observed to be strongly dependent on frequency especially for T < 413 K. This dependence of conductivity on frequency can be explained by the predominance of the hopping mechanism [23]. Hopping electrons possess a wide range of relaxation times and as a result the conductivity should increase with increasing frequency.

For higher temperatures T > 413 K, conductivity is almost constant for a wide frequency range up to a threshold frequency at which it starts increasing quite considerably. At higher temperatures and low frequency region, the conduction mechanism is mainly of the band type [24].

TABLE II

Activation energy values for Au/CuPc/Al devices at different frequencies.

| Frequency | Activation energy [eV] | | | |
|-----------|------------------------|-------|-------|--|
| [Hz] | E_1 | E_2 | E_3 | |
| 10^{2} | 0.89 | 0.43 | 0.14 | |
| 10^{3} | 0.85 | 0.36 | 0.12 | |
| 10^{4} | 0.62 | 0.24 | 0.1 | |
| 10^{5} | 0.04 | 0.16 | 0.05 | |

The conductivity versus 1000/T is illustrated in Fig. 6 for different frequencies. Activation energy is calculated

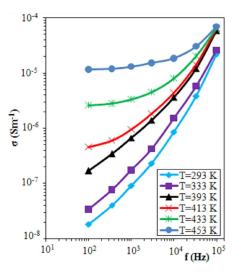


Fig. 5. Variation of AC conductivity with frequency at different temperatures.

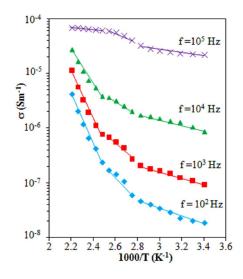


Fig. 6. Variation of AC conductivity with temperature at different frequencies.

from the slope of the plots for three different temperature regions and the obtained results are enumerated in Table II. E_1 is the activation energy in the higher temperature region and E_2 and E_3 are those in the lower temperature region. E_1 arises from the intrinsic charge carriers when the conduction mechanism is free band type, while E_2 and E_3 depend on the extrinsic conduction due to impurity scattering. In the extrinsic conduction region the charge carriers move by hopping along with the ions or electrons [25].

3.3.2. Oxygen exposure and annealing effects on conductivity

Figure 7 displays a logarithmic plot of conductivity versus frequency for samples exposed to oxygen gas with the pressure of 0.5 mbar for various exposure times of 24, 48, 168, and 336 h at room temperature. It is evident that in the low frequency range, conductivity is strongly dependent on oxygen exposure time. However, it is observed that for frequencies higher than 10^4 Hz this dependence no longer exists and all samples exhibit quite the same behavior.

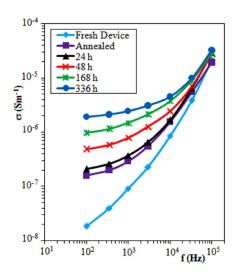


Fig. 7. Variation of AC conductivity with frequency at different exposure times to O_2 and after annealing.

The mechanism of oxidation and reduction is an important factor in the increase or decrease in conductivity. The coordination of O_2 to MPc at the air-Pc interface results in the formation of MPc⁺ and O_2^- and causes the injection of hole charge carriers into the solid [26] and hole concentration near the surface increases. Thus a rise in electrical conductivity is achieved through a doping mechanism [27] at low frequency region. Most Pcs are *p*-type semiconductors and their conductivity is enhanced after exposure to acceptor gases like O_2 [10].

The samples which were previously exposed to oxygen for 336 h went through heat treatment under high vacuum at T = 452 K for 1 h and then were cooled down to the room temperature prior to AC electrical characterization. The conductivity versus frequency for fresh devices at room temperature and the annealed samples is also plotted in Fig. 7. It is observed that conductivity of the sample decreases significantly by annealing, but it is still higher than the value obtained for the fresh devices since the complete desorption of O₂ molecules from the CuPc film was not achieved. It is probable that if the annealing temperature was higher and/or the annealing process duration was longer, complete removal of oxygen would occur and the same curve as the fresh device would appear for annealed samples.

3.3.3. Frequency and temperature effects on capacitance

The frequency dependence of capacitance in the temperature range of 293–453 K is shown in Fig. 8. At relatively high temperatures (T > 413 K) and low frequencies ($f < 10^4$ Hz), capacitance seems to be strongly dependent on frequency, but at lower temperature and higher frequency range the curves are practically parallel to the frequency axis.

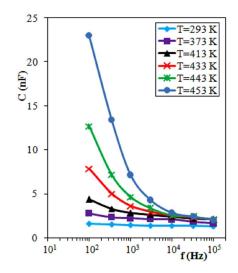


Fig. 8. Variation of capacitance with frequency at different temperatures.

According to Goswami and Goswami model [28] if our capacitor system is assumed to comprise (i) an inherent capacity element (C) unaffected by frequency (f) and temperature, (ii) a discrete resistance element (R) due to the dielectric film in parallel with C and (iii) a series resistance (r) due to lead length, etc., and if the equivalent series circuit ($C_{\rm s}$ and $R_{\rm s}$) is taken into account, the equivalent series capacitance (with $\omega = 2\pi f$) can be written as

$$C_{\rm s} = \frac{1 + \omega^2 R^2 C^2}{\omega^2 R^2 C} = \frac{1}{\omega^2 R^2 C} + C.$$
(4)

If, however, $1/\omega^2 R^2 C \ll C$, then $C_{\rm s} = C$. This condition can be satisfied in a capacitor system either by increasing R (i.e. by lowering of the temperature) or by raising ω . A suitable combination of these factors will lead to the frequency independent capacitance [28].

The capacitance as a function of temperature is plotted in Fig. 9. It can be clearly seen that at high temperatures and for frequencies less than 10^4 Hz, capacitance increases considerably with increasing temperature particularly for $f = 10^2$ Hz.

3.3.4. Oxygen exposure and annealing effects on capacitance

The variation of capacitance with frequency for samples exposed to oxygen for various exposure times is plotted in Fig. 10. It is seen that in low frequency region capacitance changes remarkably with exposure time. The increase of $C_{\rm s}$ at low frequencies can be ascribed to the increase of the conductivity of CuPc layer induced by oxygen adsorption which leads to a fall in the value of R in Eq. (4) and thus causes the measured capacitance to rise.

For frequency values higher than 10^4 Hz, however, all samples show quite the same behavior and capacitance is almost independent of frequency and reaches a minimum

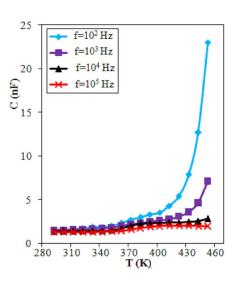


Fig. 9. Variation of capacitance with temperature at different frequencies.

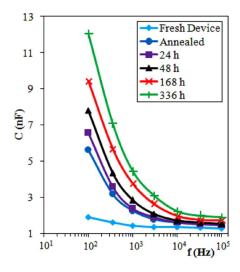


Fig. 10. Variation of capacitance with frequency at different exposure times to O_2 and after annealing.

value which is nearly constant for all samples. This effect has been studied in the previous section.

Annealing (at T = 452 K for 1 h) of the samples which were already exposed to O₂ for 336 h has greatly reduced the capacitance of the samples. This effect can be attributed to desorption of oxygen molecules from the CuPc layer. It is clearly seen in Fig. 10 that annealing has not completely reversed the effect of oxygen adsorption. The difficult process of desorption of oxygen molecules from CuPc thin film is possibly responsible for this effect. We expect full reversibility for a longer annealing duration.

3.3.5. Frequency and temperature effects on loss factor

The loss factor $(\tan \delta)$ as a function of frequency and temperature is plotted in Figs. 11 and 12. As shown in Fig. 11, loss factor is strongly dependent on frequency particularly at relatively high temperatures and low frequencies. It decreases sharply with increasing frequency to a minimum value at about $f = 10^4$ Hz and increases gradually afterwards. The behavior observed here is in good agreement with Goswami and Goswami model. Referring to this model again, loss factor is proposed as [28]:

$$\tan \delta = \frac{1}{\omega RC} + \frac{r}{\omega R^2 C} + \omega r C.$$
(5)

When $\omega R^2 C \gg r$ or $r/R \ll 1$ which is true for all cases, Eq. (5) is written as

$$\tan \delta = \frac{1}{\omega RC} + \omega rC. \tag{6}$$

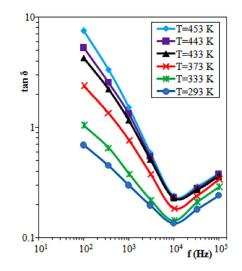


Fig. 11. Variation of loss factor with frequency at different temperatures.

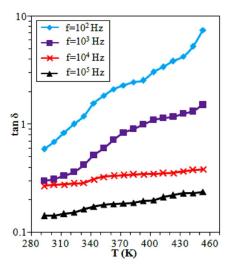


Fig. 12. Variation of loss factor with temperature at different frequencies.

Equation (6) predicts that at high frequency region the ω term is dominant whereas at low frequency region the term containing $1/\omega$ is dominant. The angular frequency

at which loss factor reaches the least value is calculated to be [28]:

$$\omega_{\min} = \left(\frac{1}{rRC^2}\right)^{1/2}.$$
(7)

It is seen in Fig. 12 that there is a gentle increase in loss factor as the temperature increases for $f \ge 10^4$ Hz, however a quite sharp increase in loss factor is observed for $f < 10^4$ Hz.

4. Conclusions

AC measurements of sandwich devices in the form of Al/CuPc/Au were carried out in temperature range of 293–453 K and frequency range of 10^2 – 10^5 Hz. FESEM images and XRD micrographs were studied to investigate the morphology of the samples. The XRD micrograph indicates the configuration of α -CuPc with the (510) plane as the preferred orientation. Optical band-gap energy of CuPc thin film was determined to be 2.81 ± 0.01 eV. At relatively high temperatures (T > 413 K) and low frequencies ($f < 10^4$ Hz), capacitance seems to be strongly dependent on frequency, but at lower temperature and higher frequency range it is generally independent of frequency. Loss factor decreased considerably with increasing frequency to a minimum value at about $f = 10^4$ Hz and increased afterwards. The activation energy decreases with increasing frequency and has a larger value at higher temperatures. The conductivity increases quite noticeably with increasing frequency particularly for T < 413 K. Exposure to oxygen was found to have a profound effect on the conductivity of samples. It is evident that in the low frequency range, conductivity is strongly dependent on oxygen exposure time. Although annealing decreased the conductivity of the samples significantly, full reversibility was not obtained. It is probable that prolonged annealing and/or higher annealing temperature can totally reverse the effects of oxygen adsorption. The AC electrical characteristics are in good agreement with Goswami and Goswami model. At low temperatures conduction mechanism is dominated by hopping-type process, whereas at high temperatures conduction is more likely to be via a band-type transport process.

References

- C.J. Liu, J.J. Shih, Y.H. Ju, Sens. Actuators B 99, 344 (2004).
- [2] T. Basova, A. Tsargorodskaya, A. Nabok, A.K. Hassan, A.G. Gurek, G. Gumus, V. Ahsen, *Mater. Sci. Eng. C* 29, 814 (2009).
- [3] A.K. Abass, R.A. Collins, A. Krier, J. Phys. Chem. Solids 54, 375 (1993).

- [4] W. Chao, X. Zhang, C. Xiao, D. Liang, Y. Wang, J. Colloid Interface Sci. 325, 198 (2008).
- [5] J. Wang, W.J. Blau, Appl. Phys. B 91, 521 (2008).
- [6] E.M. El-Menyawy, H.M. Zeyada, M.M. El-Nahass, *Solid State Sci.* **12**, 2182 (2010).
- [7] F. Yakuphanoglu, M. Arslan, Solid State Commun. 132, 229 (2004).
- [8] A. Krier, M.E. Azim-Araghi, J. Phys. Chem. Solids 58, 711 (1997).
- [9] M.E. Azim-Araghi, M.J. Jafari, Eur. Phys. J. Appl. Phys. 52, 10402 (2010).
- [10] M.E. Azim-Araghi, E. Karimi-Kerdabadi, M.J. Jafari, Eur. Phys. J. Appl. Phys. 55, 30203 (2011).
- [11] M.M. El-Nahass, F.S. Bahabri, A.A. Al-Ghamdi, S.R. Al-Harbi, *Egypt. J. Sol.* 25, 307321 (2002).
- [12] H.S. Lee, M.W. Cheon, Y.P. Park, *Trans. Electr. Electron. Mater.* **12**, 40 (2011).
- [13] M.I. Newton, T.K.H. Starke, G. McHale, M.R. Willis, *Thin Solid Films* 360, 10 (2000).
- [14] Y.L. Lee, C.Y. Hsiao, R.H. Hsiao, *Thin Solid Films* 468, 280 (2004).
- [15] A. Oprea, U. Weimar, E. Simon, M. Fleischer, H.P. Frerichs, Ch. Wilbertz, M. Lehmann, *Sens. Actuators B* 118, 249 (2006).
- [16] A.K. Hassan, R.D. Gould, J. Phys., Condens. Matter 1, 6679 (1989).
- [17] M.T. Robinson, G.E. Klein, J. Am. Chem. Soc. 74, 6294 (1952).
- [18] B.B. He, Two-dimensional X-ray Diffraction, Wiley Hoboken, NJ 2009, p. 191.
- [19] A.K. Hassan, R.D. Gould, Phys. Status Solidi A 132, 91 (1992).
- [20] M.J. Stillman, T.N. Nyokong, in: *Phthalocya-nines: Properties and Applications*, Vol. 1, Eds. C.C. Lenznoff, A.B.P. Lever, VCH, New York 1989, Ch 3, p. 133.
- [21] S.B. Brown, Introduction to Spectroscopy for Biochemistry, Academic Press, London 1980.
- [22] S. Mathew, C. Sudarsanakumar, C.S. Menon, Optoelectron. Adv. Mater.-Rapid Commun. 1, 614 (2007).
- [23] M.E. Azim-Araghi, D. Campbell, A. Krier, R.A. Collins, *Semicond. Sci. Technol.* 11, 39 (1996).
- [24] M.E. Azim-Araghi, Indian J. Pure Appl. Phys. 45, 40 (2007).
- [25] P.J Harrop, D.S. Campbell, in: Handbook of Thin Film Technology, Eds. L.I. Maissel, R. Glang, Mc-Graw Hill, New York 1983, Ch. 16, p. 16.
- [26] H.R. Kerp, K.T. Westerdui, A.T. van Veen, E.E. van Faasseu, J. Mater. Res. 16, 503 (2001).
- [27] J.D. Wright, Prog. Surf. Sci. 31, 1 (1989).
- [28] A. Goswami, A.P. Goswami, Thin Solid Films 16, 175 (1973).