Electronic Structure of CePtIn and LaPtIn Compounds

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The electronic structure of the ternary RPtIn (R = La, Ce) compounds, which crystallize in the hexagonal ZrNiAl-type structure, was studied by X-ray photoelectron spectroscopy measurements and calculation using the *ab initio* methods (linear muffin-tin orbital in the atomic sphere approximation, full potential linear muffin-tin orbital, full potential linear orbital). The results showed that the valence band in these compounds is formed by the Pt 5d and In 5s and 5p states. The band calculations with spin-orbit coupling have shown that the Ce 4f peaks consist of two peaks above the Fermi level that correspond to the Ce $4f_{7/2}$ and $4f_{5/2}$ doublet and wide peaks corresponding to the La 4f states. The analysis of Ce 3d spectra on the basis of the Gunnarsson-Schönhammer model gives hybridization of 4f orbitals with the conduction electron band equal to 170 meV.

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1. Introduction

The equiatomic ternary compounds CePtIn and LaPtIn crystallize in the hexagonal ZrNiAl structure (space group $P\bar{6}2m$) [1]. CePtIn is classified as dense Kondo material [2]. Heat capacity (C_p) measurements of CePtIn yield $C_p/T \approx 1 \text{ J}/(\text{K}^2 \text{ mol})$ at 60 mK that indicates heavy fermion properties. The specific heat and electrical resistivity studies indicate the absence of magnetic order up to 50 mK [3]. Temperature dependence of the electrical resistivity of LaPtIn indicates the metallic character and results for (Ce_{1-x}La_x)PtIn system indicate its change of the properties with increasing La content [4].

In the CeTIn compounds with increase in the unit cell volume the magnetic properties change from the Pauli paramagnet for CeNiIn to antiferromagnetical order found in CePdIn ($T_{\rm N} = 1.8$ K) and CeAuIn ($T_{\rm N} =$ 5.7 K). Simultaneously, a decrease in the Kondo temperature from $T_{\rm K} = 94$ K for CeNiIn to 3.3 K for CePdIn was reported [5]. The CePtIn compound with $T_{\rm K} = 11$ K and no magnetic order lie between those two border regions. Determination of the parameters of the electronic structure of this compound is interesting for explanation of the described changes in physical properties.

The photoemission method (PES) is considered to be the most successful way to study the nature of the Ce 4fstates [6]. The main aim of our investigations was to find correlations between the data on the electronic structure of the investigated compounds and hybridization of the 4f level in the photoemission spectra of the valence band (VB). In practice, the small cross-section of the 4f states comparing to the d ones makes the analyses very difficult when using classical photoemission. More reliable information about the 4f states can be achieved using resonant photoemission. Also the Ce 3d core-level photoemission lines have been widely used to diagnose hybridization strength of the 4f levels in the Ce-based systems [6].

In this work we present the data concerning the electronic structure of RPtIn (R = La, Ce) compounds supported by *ab initio* calculations and PES results.

2. Experimental details

Polycrystalline samples of RPtIn (R = La, Ce) compounds were obtained by arc melting of stoichiometric amounts of the elements (Merck AG: Ce and La of 99.9% purity, Pd and In of 99.99% purity) under a high-purity argon atmosphere on a water-cooled copper hearth, using a titanium-zirconium alloy as a getter. In order to ensure good homogeneity the buttons were turned over and remelted several times. The final weight losses were lower than 0.1%. Subsequently the samples were annealed in evacuated quartz tubes at 800 °C for one week. The quality of the products was checked by X-ray powder diffraction using Philips PE-3710 apparatus (Cu K_{α} radiation) at room temperature. All of the reflections were indexed within the hexagonal structure of ZrNiAl-type. For processing the collected data the Rietveld-type program Fullprof was used [7]. The determined structure parameters for RPtIn (R = La, Ce) compounds are in good agreement with the previous data [1].

The X-ray photoelectron spectra (XPS) were obtained at room temperature using a Multipurpose Electron Spectrometer PHI5700/660 with monochromatised radiation Al K_{α} ($h\nu = 1486.6$ eV) radiation source and the overall resolution of 0.3 eV. Binding energies were referred to the Fermi level ($E_{\rm F} = 0$). The spectrometer was calibrated using the Cu $2p_{3/2}$ (932.5 eV), Ag $3d_{5/2}$ (368.1 eV) and Au $4f_{7/2}$ (84.0 eV) core-level photoemission spectra. Surface of the samples was mechanically cleaned by scraping with a diamond file in a preparation chamber under high vacuum conditions (10^{-9} mbar). The Shirley method was used to subtract background; the experimental spectra obtained in this manner were numerically refined using the 80% Gaussian and 20% Lorentzian peak shape [8].

TABLE I

Structural parameters of CePtIn and LaPtIn at room temperature.

Compound	LaPtIn	CePtIn
a [Å]	7.823(1)	7.706(1)
c [Å]	4.165(1)	4.097(1)
a/c	1.878(1)	1.881(1)
V [Å ³]	220.74(11)	210.69(10)
R 3g $(x, 0, 1/2)$	0.586(1)	0.0.588(1)
In 3f $(x, 0, 0)$	0.249(1)	0.251(1)
Pt1 1b $(0, 0, 1/2)$		
${\rm Pt}2 2{\rm c} (1/3, 2/3, 0)$		

The electronic density of states (DOS) of RPtIn (R = La, Ce) has been calculated by the spin-polarized selfconsistent tight binding linear muffin-tin orbital (TB LMTO) method in the atomic sphere approximation (ASA) [9, 10] using the structural parameters from Table I. The scalar relativistic approximation for band electrons and the full relativistic treatment of frozen core electrons were applied. The exchange correlation potential was assumed in the form of von Barth and Hedin [11] and gradient corrections were also included [12].

3. Results

The calculated band structure of CePtIn is shown in Fig. 1. Above the Fermi level (0-0.5 eV) the Ce 4f state is observed. In the region of 3-4 eV the Pt 5d states is detected. On the basis of this data the density of states was calculated. The contribution from Ce or La, Pt and In and the total density of states for CePtIn and LaPtIn calculated by scalar relativistic method are shown in Fig. 2.

Both methods of calculation give nearly similar DOS. Above the Fermi level ($E_{\rm F} = 0$ eV) the doublet corresponding to the Ce $4f_{3/2}$ and Ce $4f_{5/2}$ is observed. In the 3–4.5 eV regions the states of Pt 5*d* are detected. At 6 eV the peaks of In 5*s* are observed whereas In 5*p* states form a band between 0 and 4.5 eV. Similar DOS is observed for LaPtIn however the La 4*f* states form the wide peak at 2.5 eV above the Fermi level.



Fig. 1. Band structure of CePtIn.



Fig. 2. Total density of states and the contribution from Ce or La, Pt and In to the total density of states for CePtIn (a) and LaPtIn (b) calculated in the scalar--relativistic approximation.

TABLE II Atomic concentration in RPtIn (R = La, Ce) compounds.

Compound	R $3d_{5/2}$	Pt $4f_{7/2}$	In $3d_{5/2}$
CePtIn	30.70	33.4	35.90
${ m LaPtIn}$	34.03	32.12	33.85

The calculated values of the density of states at the Fermi level $N(E_{\rm F})$ for both methods of calculation are shown in Table II.

TABLE III

Density of states at the Fermi level $N(E_{\rm F})$ for both methods of calculations.

	Scalar realistic	Full relativistic
LaPtIn	169.9	$157.8 \ states$
${\rm CePtIn}$	313.2	$350.1 \ { m states}$

These results indicate that the increase in $N(E_{\rm F})$ for CePtIn is caused by the Ce 4f states. The calculated occupation of the Ce 4f level is equal to 0.97.

Figure 3 shows the XPS spectra of CePtIn and LaPtIn collected in a wide 0-1400 eV binding energy range, the binding energies are related to the Fermi level ($E_{\rm F}$ = 0 eV). The results of the analysis of the intensity of the selected peaks corresponding to R $3d_{5/2}$, Pt $4f_{7/2}$ and In $3d_{5/2}$ states are given in Table III. The obtained results indicate that in both compounds the ratio 1:1:1 is retained. In general, the most important data for CePtIn were obtained from three regions: VB (binding energy 0-10 eV), Ce 4d states (100-120 eV) and Ce 3d states (880–930 eV). The XPS VB for both compounds is presented in Fig. 4. These spectra present mainly the dominant contribution of the Pt states. For CePtIn this corresponds to two peaks: at 3.2 and 4.5 K. Similar two peaks at 3.2 and 4.4 eV are observed in LaPtIn. The adequate values of spin-orbit splitting are equal to 2.5 and 1.2 eV.



Fig. 3. XPS Al K_{α} spectra of LaPtIn and CePtIn collected at room temperature within 0–1400 eV binding energy range (Al K_{α} radiation).

A broad peak at 6.1 eV for CePtIn and at 6.6 eV for LaPtIn corresponds to In 5s states. Near the Fermi level two peaks at 0.5 and 2.1 eV for CePtIn and one at 0.5 and 1.9 eV for LaPtIn were observed. The peaks at 0.5 eV are connected to the $5d6s^2$ state where the peak at 1.9 eV for LaPtIn is connected with the In 5p state. For CePtIn the peak at 2.1 eV is probably connected with the Ce 4f and In 5p state.

At about 110 eV a broad peak of the Ce 4f states can be noticed (see Fig. 5). The Ce $4d_{5/2}$ lines is decomposed into four peaks with the binding energies equal to



Fig. 4. Deconvoluted photoemission valence band spectra for LaPtIn and CePtIn.



Fig. 5. Deconvoluted Ce 4d spectrum in CePtIn and La 4d spectrum in LaPtIn.

107.9, 110.0, 112.2, and 115.0 eV. Unfortunately analysis of the structure of this peak is possible. In LaPtIn in this region two peaks corresponding to the La $4d_{5/2}$ and La $4d_{3/2}$ are observed at 102.55 and 105.85 eV and the spin-orbit splitting is equal to $\Delta = 3.3$ eV. These data are in good agreement with the data for the pure elements 102.5, 105.3, and 2.8 eV [13].

Figure 6 illustrates the Ce 3d XPS spectrum of the CePtIn compound. Two lines corresponding to the spin-

orbit splitting equal to 18.4 eV ($3d_{3/2}$ and $3d_{5/2}$ doublets) are observed. Additionally, each line consists of two sublines corresponding to the $3d^94f^1$ and $3d^84f^2$ states. The analysis of these components was based on the Doniach– Šunjić theory [14]. Within this approach the intensity ratio $r_1 = I(f^2)/[I(f^2) + I(f^1)]$ can be calculated ($I(f^n)$ is the integral intensity of relevant components). From r_1 value it is possible to determine the coupling parameter Δ basing on the calculations of the intensity ratio r_1 as a function of Δ previously published [6]. The coupling parameter Δ is defined as $\pi V^2 \rho_{\max}$, where V is the hybridization and ρ_{\max} is the maximum of the density of conduction electron states. The value of the intensity ratio r_1 is equal to 0.33 which corresponds to the Δ parameter of 170 eV.



Fig. 6. Deconvoluted X-ray photoelectron spectra of Ce $3d_{5/2}$ and Ce $3d_{3/2}$ electron states in CePtIn.

TABLE IV

The binding energies and the spin-orbit splitting Δ_{S-O} (in eV) in RPtIn (R = La, Ce) compounds, compared with the data for pure element.

	LaPtIn	CePtIn	Pure element [13]
In $4d_{5/2}$	16.8	16.85	16.6
$4d_{3/2}$	17.7	17.81	17.4
$\Delta_{\mathrm{S-O}}$	0.9	1.0	0.8
Pt $4f_{7/2}$	71.6	71.6	71.1
$4f_{5/2}$	74.9	74.9	74.7
$\Delta_{\rm S-O}$	3.3	3.3	3.6
Pt $4d_{5/2}$	314.8	314.8	314.6
$4d_{3/2}$	331.6	331.6	331.6
$\Delta_{\mathrm{S-O}}$	16.8	16.8	17.0
In $3d_{5/2}$	443.7	443.8	443.9
$3d_{3/2}$	451.3	451.3	451.4
$\Delta_{\rm S-O}$	7.6	7.5	7.5

The evidence of the intermediate valence of Ce ion is not clear because the Ce $3d^94f^0$ peak at 916.2 eV and the Ce 4d one connected with the f^0 final state near 122.5 eV are broad.

4. Summary

In this work we investigated the electronic structure of CePtIn and LaPtIn. The small difference in the valence band of these compounds was observed. From the 3d spectrum the hybridization parameter equal to 170 eV was determined which is close to these value reported for the isostructural CePdIn (130 eV) [5], in spite that these compounds have different properties. The presented data do not give explanation of anomalous properties of CePtIn. The Ce 3d and Ce 4d spectra do not give any evidence of intermediate valence of Ce³⁺ in CePtIn.

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