

---

Proceedings of the XIII National School of Superconductivity, Łądek Zdrój 2007

## Thermoelectric Power of MgB<sub>2</sub> Single Crystals Doped with Holes and Electrons

K. OGANISIAN<sup>a</sup>, K. ROGACKI<sup>a</sup>, C. SULKOWSKI<sup>a</sup>,  
N.D. ZHIGADLO<sup>b</sup>, S. KATRYCH<sup>b</sup> AND J. KARPINSKI<sup>b</sup>

<sup>a</sup>Institute of Low Temperature and Structure Research, Polish Academy of Sciences, P.O. Box 1410, 50-950 Wrocław, Poland

<sup>b</sup>Laboratory for Solid State Physics, ETH, 8093 Zurich, Switzerland

The thermoelectric power of MgB<sub>2</sub> single crystals, non-substituted, single-substituted with C, and co-substituted with C and Li, was investigated in the temperature range from 10 to 300 K. Both the in-plane ( $S_{ab}$ ) and the out-of-plane ( $S_c$ ) thermoelectric powers are positive for non-substituted crystals and both  $S_{ab}$  and  $S_c$  change a sign for crystals doped with electrons when C is substituted for B in the amount larger than 6 at.%. The substitution of C, which supplies electrons mainly into the  $\sigma$  band, reduces the thermoelectric power anisotropy, most likely by increasing the interband scattering. When Li is additionally co-substituted for Mg,  $S_{ab}$  remains unchanged but  $S_c$  increases substantially. The Li substitution donates holes into the  $\pi$  band rather than into the  $\sigma$  band and in this way modifies the  $\pi$  band properties, including the possible increase in the intraband scattering.

PACS numbers: 74.25.Fy, 74.70.Ad, 74.62.Dh

### 1. Introduction

Structure and electronic properties of MgB<sub>2</sub> can be modified in a controlled way by substitutional chemistry. The superconducting transition temperature,  $T_c$ , upper critical field and its anisotropy, gaps width, intra- and interband scattering, defect structure, etc., can be tuned in the wide range if the doping with holes and electrons is possible. Introducing Li, Al, Mn, Fe, and C into the melt, one can grow MgB<sub>2</sub> crystals where Mg and B are partially substituted and thus the crystals are doped with holes (Li) and electrons (Al, Fe, C) or they are substituted isovalently (Mn) [1–7]. The substituted crystals can be used to study the influence of doping and/or intra- and interband scattering on the anisotropic normal and superconducting properties of MgB<sub>2</sub>. In this work we examine the temperature behavior of the thermoelectric power of MgB<sub>2</sub> single crystals with various substitutions. Both electron and hole doped crystals have been investigated, including crystals co-substituted with C–Li to study the possible effect of electron-hole compensation [7].

## 2. Experimental

Single crystals of C, Li, and C–Li substituted  $\text{MgB}_2$  were grown under high pressure using a cubic anvil press. The applied temperature and pressure conditions for the growth of C substituted crystals were determined in our previous studies [3–5]. The successful method to obtain Li and C–Li substituted crystals was described in Ref. [7]. The C content in the  $\text{MgB}_2$  crystals has been estimated from changes in the  $a$  lattice parameter. The amount of Li has been obtained from the structure refinement as described in Ref. [7]. Using the high pressure synthesis we have grown single crystals with dimensions up to  $1.5 \times 1 \times 0.1 \text{ mm}^3$  (see Fig. 1).

The thermoelectric power,  $S$ , has been measured for the temperature gradient,  $\Delta T$ , applied both in the  $ab$  plane,  $S_{ab}$ , and along the  $c$  axis,  $S_c$ , of the  $\text{MgB}_2$  crystal. The crystal was placed between two Cu blocks fitted with two heaters and two precisely calibrated Pt thermometers. The heaters kept a certain temperature difference between the Cu blocks and the potential difference between the top and the bottom of the sample was measured by the sensitive voltmeter. For  $\Delta T$  from about 2 to 4 K, the thermoelectric power was gradient independent and the value of  $S$  for this gradient has been taken as a correct result. The measuring setup was tested with Sb and Pb standards and with a reference crystal of  $\text{YBa}_2\text{Cu}_3\text{O}_7$ . An error of the absolute value of  $S$  was estimated to be typically below  $0.2 \mu\text{V}/\text{K}$ .

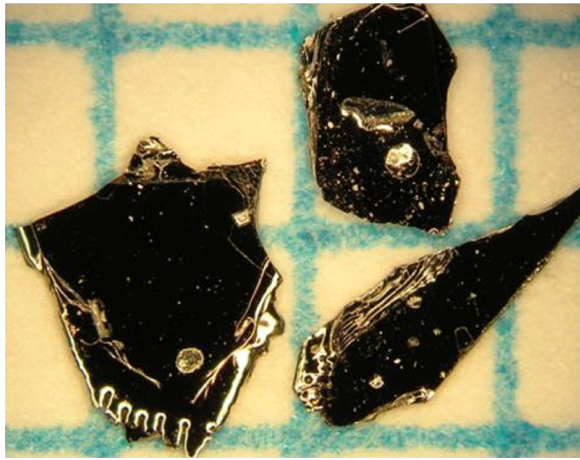


Fig. 1.  $\text{MgB}_2$  single crystals (with various substitutions) grown in the high pressure and high temperature cubic anvil press. Scale is 1 mm.

## 3. Results and discussion

Single crystals of  $\text{Mg}(\text{B}_{1-x}\text{C}_x)_2$  with C content changing from 2% to 7% have been studied as well as the crystals which were non-substituted, Li substituted, and C–Li co-substituted, for comparison. The influence of the C substitution on the

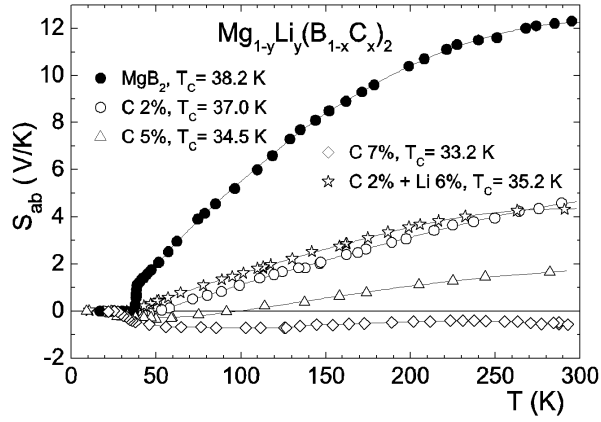


Fig. 2. Temperature dependence of the thermoelectric power  $S_{ab}$  of  $Mg_{1-y}Li_y(B_{1-x}C_x)_2$  single crystals; non-substituted, carbon-substituted with  $x = 0.02, 0.05,$  and  $0.07,$  and C–Li co-substituted with  $x = 0.02$  and  $y = 0.06.$   $S_{ab}$  is measured with the temperature gradient in the  $ab$  plane.

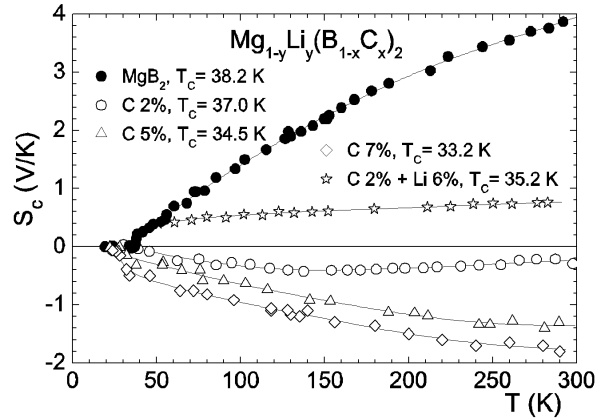


Fig. 3. Temperature dependence of the thermoelectric power  $S_c$  of  $Mg_{1-y}Li_y(B_{1-x}C_x)_2$  single crystals; non-substituted, carbon-substituted with  $x = 0.02, 0.05,$  and  $0.07,$  and C–Li co-substituted with  $x = 0.02$  and  $y = 0.06.$   $S_c$  is measured with the temperature gradient along the  $c$  axis.

thermoelectric power is shown in Figs. 2 and 3, for measurements performed with a temperature gradient in the  $ab$  plane and along the  $c$  axis, respectively. For the non-substituted  $MgB_2,$   $S_{ab}$  decreases smoothly with decreasing temperature down to  $T_c = 38.2$  K, then  $S_{ab}$  drops sharply to zero, as expected. Similar temperature dependence has been observed for  $S_c,$  however the absolute value of  $S_c$  is much lower. The thermoelectric power anisotropy, defined as  $S_{ab}/S_c,$  changes from 3 to 5 when temperature decreases from 300 K to  $T_c.$  The raise of the anisotropy with decreasing temperature means that  $S_c$  drops faster than  $S_{ab}$  suggestive of

the interband-scattering reduction. When C is substituted for B, both  $S_{ab}$  and  $S_c$  become smaller with the C content and, for a larger amount of C, both of them change the sign to negative. For  $S_{ab}$ , the substitution of 7% of C is necessary to observe the negative values. For  $S_c$ , however, the only 2% substitution of C is enough to modify the Fermi surface, change the charge carriers distribution significantly, and in this way change the sign of  $S$ . Thus, the substitution of C for B, which donates electrons mostly into the 2D  $\sigma$  band, strongly influences the electronic structure and results in substantial changes of the thermoelectric properties in both directions, parallel and perpendicular to the  $ab$  plane. The difference between  $S_{ab}$  and  $S_c$  decreases with the C content, so we conclude that the C substitution not only alters the  $\sigma$  band filling but probably also increases the interband scattering noticeably. Both  $S_{ab}$  and  $S_c$  saturate with the C content and the extrapolated saturation values of  $S_{ab} \approx -2 \mu\text{V/K}$  and  $S_c \approx -3 \mu\text{V/K}$  have been observed at 300 K.

As we mentioned above, Li substitutes for Mg and dopes  $\text{MgB}_2$  with holes. Thus, an intriguing question arises if it is possible to compensate the effect of C substitution, which dopes  $\text{MgB}_2$  with electrons, by the C–Li co-substitution. For this reason, the thermoelectric power of the C–Li co-substituted crystals has been also studied and the results are included in Figs. 2 and 3. The substitution of 2% of C reduces  $T_c$  slightly from 38.2 K to 37 K. The co-substituted crystal with 2% of C and 6% of Li shows additional reduction of  $T_c$  to 35.2 K. However, the thermoelectric power  $S_{ab}$  for crystals substituted with 2% of C and co-substituted with 2% of C and 6% of Li is practically the same. Thus, Li does not influence the thermoelectric properties of the  $ab$  plane, that can be explained assuming that holes coming from Li occupy almost entirely the  $\pi$  band. This is supported by the results of first-principles calculations, which show that only a small amount of the Li-introduced holes may dope the  $\sigma$  band [8]. On the other hand,  $S_c$  for the crystal co-substituted with 2% of C and 6% of Li is distinctly larger than  $S_c$  for the crystal single-substituted with 2% of C. Thus the anisotropy of  $S$  for the C–Li co-substituted crystal is smaller than for the C single-substituted one. Assuming that Li donates holes mainly into the  $\pi$  band, this feature may be explained as a result of changes in the electronic properties of this 3D band only, including the intraband scattering. The interband scattering, however, seems not to be influenced by the Li co-substitution (up to 6%) in any substantial degree. This speculation has to be verified by additional experiments with crystals single-substituted with various amount of Li, which are under way. Preliminary results show that 6% of Li single-substituted for Mg changes (if any) both  $T_c$  and  $S$  only slightly.

#### 4. Conclusions

In conclusion, availability of sizeable  $MgB_2$  crystals doped with electrons and holes enabled us to perform advanced studies of the anisotropic thermoelectric properties of the two-band superconductor. The thermoelectric power is strongly influenced by doping when C is substituted for B. This substitution, which introduces electrons into the  $\sigma$  band, results in a substantial decrease in both  $S_{ab}$  and  $S_c$ , leading to the change of their sign to negative for the C content equal to about 7 and 2%, respectively. The observed decrease in  $S$  indicates more metallic and electron-conducting character of the C substituted crystals. The substitution of C reduces the thermoelectric power anisotropy, most likely by increasing the interband scattering. When 6% of Li is additionally introduced into the crystal with 2% of substituted C,  $S_{ab}$  practically does not change, however  $S_c$  increases. This effect suggests that holes, which dope most likely the  $\pi$  band, change the conducting features of this band only. The electron- and hole-doping effects seem to be asymmetric for the thermoelectric normal-state properties, similarly as observed for the superconducting properties [7]. In order to fully understand the relation between doping and the intra- and interband scattering, additional experiments are necessary with crystals single-substituted with Li and Al, and crystals co-substituted with Li–Al. Such crystals have been recently successfully grown.

#### Acknowledgments

This work was supported by the Polish Ministry of Science and Higher Education under a research project No. N202 131 31/2223 for the years 2006-2009 and by the Swiss National Science Foundation through NCCR pool MaNEP.

#### References

- [1] T. Masui, S. Lee, S. Tajima, *Phys. Rev. B* **70**, 024504 (2004).
- [2] S. Lee, T. Masui, A. Yamamoto, H. Uchiyama, S. Tajima, *Physica C* **397**, 7 (2003).
- [3] S.M. Kazakov, R. Puzniak, K. Rogacki, A.V. Mironov, N.D. Zhigadlo, J. Jun, Ch. Soltmann, B. Batlogg, J. Karpinski, *Phys. Rev. B* **71**, 024533 (2005).
- [4] J. Karpinski, N.D. Zhigadlo, G. Schuck, S.M. Kazakov, B. Batlogg, K. Rogacki, R. Puzniak, J. Jun, E. Müller, P. Wägli, R. Gonnelli, D. Daghero, G.A. Ummarino, V.A. Stepanov, *Phys. Rev. B* **71**, 174506 (2005).
- [5] K. Rogacki, B. Batlogg, J. Karpinski, N.D. Zhigadlo, G. Schuck, S.M. Kazakov, P. Wägli, R. Puźniak, A. Wiśniewski, F. Carbone, A. Brinkman, D. van der Marel, *Phys. Rev. B* **73**, 174520 (2006).
- [6] C. Krutzler, M. Zehetmayer, M. Eisterer, H.W. Weber, N.D. Zhigadlo, J. Karpinski, A. Wisniewski, *Phys. Rev. B* **74**, 144511 (2006).
- [7] J. Karpinski, N. Zhigadlo, S. Katrych, K. Rogacki, B. Batlogg, M. Tortello, R. Puzniak, *Phys. Rev. B* **77**, 214507 (2008).
- [8] F. Bernardini, S. Massidda, *Phys. Rev. B* **74**, 014513 (2006).