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## Specific Heat of YbNi<sub>4</sub>Si Compound

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The specific heat of YbNi<sub>4</sub>Si has been analyzed considering the electronic contribution and the lattice contributions in frames of the Debye model. Based on the specific heat measurements, the electronic specific heat coefficient  $\gamma = 25 \text{ mJ mol}^{-1} \text{ K}^{-2}$  and the Debye temperature  $\theta_D = 320 \text{ K}$  were derived. This small value shows that YbNi<sub>4</sub>Si cannot be classified as a heavy fermion system. These studies are completed by magnetic susceptibility, X-ray photoemission spectroscopy, and electrical resistivity.

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### 1. Introduction

The intermediate valence compounds pose one of the most challenging problems of strongly correlated electron systems. Different ingredients contribute to the complexity of these fascinating systems: the presence of strong Kondo interactions, the level structure of the crystal electric field split  $f$ -orbitals, a different hybridization between each level and the conduction band, and the eventual coherence effects and magnetic interactions introduced by the periodicity of the Kondo lattice.

A great interest in the Yb-based compounds comes from the fact that Yb can exhibit two valence states related to the nonmagnetic  $4f^{14}$  (Yb<sup>2+</sup>) and magnetic  $4f^{13}$  (Yb<sup>3+</sup>) electronic configurations. Due to this feature Yb is often treated as a hole counterpart of Ce, which also shows two close in energy, electronic configurations: the magnetic  $4f^1$  (Ce<sup>3+</sup>) and the nonmagnetic  $4f^0$  (Ce<sup>4+</sup>).

In our paper we analyze the heat capacity data of YbNi<sub>4</sub>Si. To verify fully the magnetic contribution due to the Yb  $f$  states the nonmagnetic reference compound YNi<sub>4</sub>Si is also measured.

### 2. Experimental

The sample preparation was described in detail in Ref. [1]. The room temperature powder X-ray diffraction pattern revealed that the studied YbNi<sub>4</sub>Si com-

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pound was single-phase. Heat capacity measurements were performed by PPMS commercial device (Quantum Design) in the temperature range of 2–300 K by the relaxation method using the two-tau model.

### 3. Results and discussion

$\text{YbNi}_4\text{Si}$  is paramagnetic and its magnetic susceptibility follows the Curie–Weiss law  $\chi(T) = C/(T - \theta)$  with  $\mu_{\text{eff}} = 4.15\mu_{\text{B}}/\text{f.u.}$  and the paramagnetic Curie temperature  $\theta \approx 0$  K [1]. This effective magnetic moment is slightly reduced with respect to the free  $\text{Yb}^{3+}$  ion ( $4.54\mu_{\text{B}}$ ) and suggests that the valence of the ytterbium ions in  $\text{YbNi}_4\text{Si}$  is close to 3+. The  $\text{Yb}^{2+}$  and  $\text{Yb}^{3+}$  peaks observed by X-ray photoemission spectroscopy (XPS) in the valence band region confirm the domination of the  $\text{Yb}^{3+}$  valence state [1] (Fig. 1).

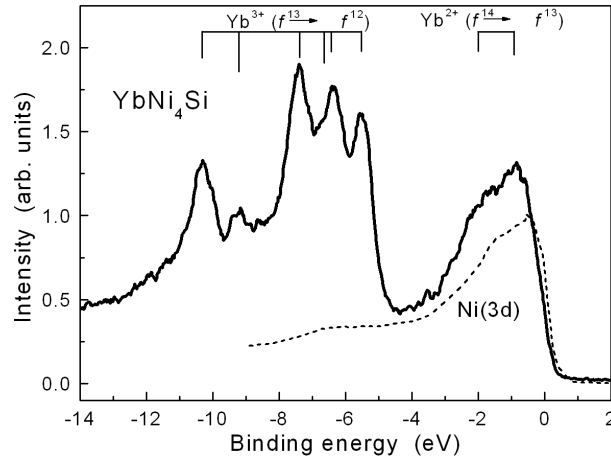


Fig. 1. Valence band region of the  $\text{YbNi}_4\text{Si}$  compound measured by X-ray photoemission spectroscopy.

The temperature dependence of the electrical resistivity  $\rho(T)$  shows the standard Fermi-liquid behavior with  $T^2$  dependence up to about 60 K.

The total specific heat consists of electronic and phonon contributions. The electronic specific heat coefficient  $\gamma$  provides information concerning the conduction band density of states at the Fermi level. Figure 2 shows the temperature dependence of the specific heat for  $\text{YNi}_4\text{Si}$  and  $\text{YbNi}_4\text{Si}$  compounds.

$\text{YNi}_4\text{Si}$  is nonmagnetic; therefore,  $C_p(T)$  can be described by the standard formula:

$$C_p(T) = \gamma T + 9NR \left( \frac{T}{\theta_D} \right)^3 \int_0^{\theta_D/T} \frac{x^4 \exp(x) dx}{[\exp(x) - 1]^2}, \quad (1)$$

where the first and the second term correspond to the electronic and the phonon contribution, respectively.  $N = 6$  is the number of the atoms in the formula unit and  $x = \hbar\omega/k_{\text{B}}T$ . From the fit to the experimental points (Fig. 2) we get

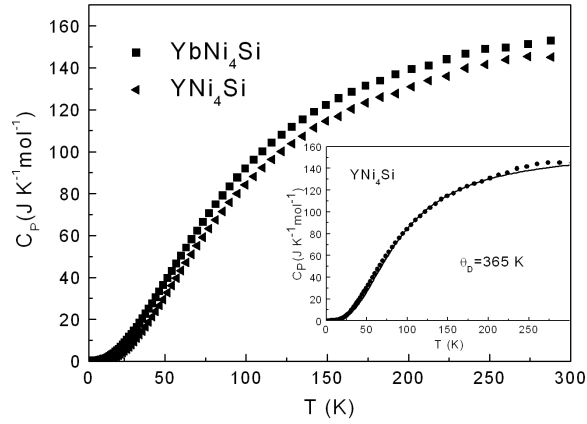


Fig. 2. Temperature dependence of the specific heat of  $\text{YbNi}_4\text{Si}$  and  $\text{YNi}_4\text{Si}$ . (Inset) Specific heat of the  $\text{YNi}_4\text{Si}$  compound fitted with Eq. (1) (solid line).

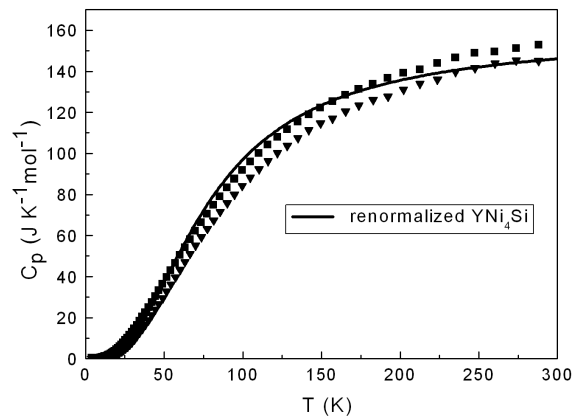


Fig. 3. The specific heat of  $\text{YbNi}_4\text{Si}$  (squares) and  $\text{YNi}_4\text{Si}$  (triangles). Solid line shows the  $\text{YNi}_4\text{Si}$  data after renormalization of the mass in respect to the Yb-based compound.

the Debye temperature  $\theta_D = 365$  K and the electronic specific heat coefficient  $\gamma = 13$   $\text{mJ mol}^{-1}\text{K}^{-2}$ . A similar analysis has been carried out for  $\text{YbNi}_4\text{Si}$  yielding  $\theta_D = 320$  K and  $\gamma = 25$   $\text{mJ mol}^{-1}\text{K}^{-2}$ , respectively. However, if a significant magnetic contribution is expected for the Yb-based compound then the simple Debye model does not describe the lattice specific heat accurately enough. In the absence of a measurable magnetic part the Debye temperature of  $\text{YbNi}_4\text{Si}$  should be easily obtained from the data for  $\text{YNi}_4\text{Si}$  by considering a correction to account for the difference between the molar masses of the  $\text{YbNi}_4\text{Si}$  and  $\text{YNi}_4\text{Si}$  compounds. The renormalization ratio [2] is expressed as:

$$\frac{\theta_D(R_m X_n Y_p)}{\theta_D(R'_m X_n Y_p)} = \left( \frac{m(M_{R'})^{3/2} + n(M_X)^{3/2} + p(M_Y)^{3/2}}{m(M_R)^{3/2} + n(M_X)^{3/2} + p(M_Y)^{3/2}} \right)^{1/3}, \quad (2)$$

where  $M_R$ ,  $M_X$  and  $M_Y$  are the molar masses of the  $R$ ,  $X$ , and  $Y$  atoms, respectively,  $\theta_D$  is the Debye temperature and can be used to evaluate the lattice contribution of the magnetic compound  $R_m X_n Y_p$  by rescaling the data for the nonmagnetic  $R'_m X_n Y_p$  reference system.

Figure 3 presents the specific heat of YbNi<sub>4</sub>Si (squares) and YNi<sub>4</sub>Si (triangles) together with the renormalized data of YNi<sub>4</sub>Si (solid line). The Debye temperatures of YbNi<sub>4</sub>Si obtained from the direct fit with Eq. (1) and from the renormalization of YNi<sub>4</sub>Si are nearly identical, which indicates that the magnetic contribution to the specific heat of YbNi<sub>4</sub>Si is very small, therefore we omit its analysis.

#### 4. Conclusions

It has been found that the YbNi<sub>4</sub>Si compound does not order magnetically down to 4 K. Nearly in the whole temperature range studied the magnetic susceptibility follows a Curie law with  $\mu_{\text{eff}} = 4.15\mu_B/\text{f.u.}$  This effective magnetic moment is close to the value expected for the  $4f^{13}$  configuration ( $4.54\mu_B$ ). The Yb<sup>2+</sup> and Yb<sup>3+</sup> peaks observed by XPS in the valence band region confirm the domination of the Yb<sup>3+</sup> valence state. The heat capacity measurements have shown that YbNi<sub>4</sub>Si is characterized by  $\gamma = 25 \text{ mJ mol}^{-1}\text{K}^{-2}$  and  $\theta_D = 320 \text{ K}$ . A mass correction to the nonmagnetic reference sample YNi<sub>4</sub>Si indicates that the magnetic contribution to the specific heat of the Yb-based compound is very small.

#### References

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- [2] M. Bouvier, P. Lethuillier, D. Schmitt, *Phys. Rev. B* **43**, 13137 (1991).