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## Thermoelectric Power of CeNi<sub>4</sub>Si and YbNi<sub>4</sub>Si Compounds

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The thermoelectric power was measured from 4.2 to 300 K for CeNi<sub>4</sub>Si and YbNi<sub>4</sub>Si. The thermoelectric power was analysed in the framework of the phenomenological resonance model. According to the model the dominant contribution to thermopower is caused by scattering between electrons of a broad *s*-band and a narrow *f*-band with the Lorentzian shape. The electron-hole analogy is reflected in the thermoelectric power behaviour of the investigated compounds.

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### 1. Introduction

The intermetallics CeNi<sub>4</sub>Si and YbNi<sub>4</sub>Si belong to large family of rare-earth RNi<sub>4</sub>Si (R = rare earth) compounds, crystallising in the hexagonal CeCu<sub>5</sub>-type structure (space group *P6/mmm*) [1]. The first compound has been reported to show paramagnetic properties down to 2 K. However, the effective moment deduced from a fit of magnetic data to the Curie-Weiss law is rather low (0.52  $\mu_B$ ), if one compares to that of the free Ce<sup>3+</sup> ions (2.54  $\mu_B$ ). This behaviour was interpreted in terms of intermediate valence of the Ce atoms and may be supported by the X-ray photoemission spectroscopy (XPS) experiments [1]. The latter have shown the *f* occupancy  $n_f$  to be 0.91 and the hybridisation energy  $\Delta = 36$  meV between *f* states and conduction electrons.

Moreover, for CeNi<sub>4</sub>Si one finds a minimum in the temperature dependence of the electrical resistivity at 8 K, resembling the behaviour of a single-ion Kondo system. YbNi<sub>4</sub>Si seems to be also an intermediate valence compound [1]. Indeed, the XPS experiments have indicated that the majority contribution to the XPS

spectra comes just from the  $\text{Yb}^{3+}$  ions. The magnetic data consistently with the XPS data show a wide temperature range of the Curie law behaviour with the effective moment ( $4.15 \mu_{\text{B}}$ ) close to that expected for the  $f^{13}$  configuration ( $4.54 \mu_{\text{B}}$ ). The information about Ce and Yb compounds mentioned above makes up new support for the concept of electron-hole analogy between Ce and Yb ions. This could be understood since the Yb atoms can appear in two valence states, nonmagnetic ( $f^{14}$ ) and magnetic ( $f^{13}$ ). In turn, Ce, which is the electron counterpart of Yb may possess magnetic ( $f^1$ ) and nonmagnetic ( $f^0$ ) valence states. Because the thermoelectric power is sensitive to the curvature of the bands at the Fermi level, the comparison of thermoelectric power (TEP) between  $\text{CeNi}_4\text{Si}$  and  $\text{YbNi}_4\text{Si}$  is highly desired. Thus, in the contribution we present experimental data of the thermoelectric power for these compounds and discuss them shortly in the framework of the existing theories.

## 2. Experimental

TEP was measured in the temperature range 4.2–300 K using the differential method. A temperature gradient of about 1 K was applied along the sample length and was determined using a pair of Au-0.07% Fe/chromel thermocouple. The absolute uncertainty of the measured Seebeck coefficient  $S$  is less than  $0.5 \mu\text{V}/\text{K}$ .

## 3. Results

The behaviour of TEP of Ce- and Yb-based intermetallics compounds often differs from those of the ordinary metals and their alloys. Instead of the magnitude of several  $\mu\text{V}/\text{K}$ , the TEP of compounds with unstable  $f$ -electron shells may attain a giant value, i.e., of the order of a hundred  $\mu\text{V}/\text{K}$ . Furthermore, in addition to the diffusion mechanism that appears as a linear temperature dependence term in the  $S(T)$  function, there are other contributions, for instance those, related to the Kondo effect, crystalline electric field or scattering of conduction electrons on a narrow  $f$  band. As a result, the DOS structure of these compounds near  $E_{\text{F}}$  is rather complicated, and it can lead to the appearance of several extrema in the  $S(T)$  curve.

Figure 1a shows the temperature dependence of the thermoelectric power for  $\text{CeNi}_4\text{Si}$ . TEP is found to be negative at room temperature. It increases steadily with decreasing temperature until reaches a maximum value of  $2.5 \mu\text{V}/\text{K}$  at  $T_{\text{max}} = 55$  K. Furthermore, the thermopower exhibits a minimum of  $\approx -1 \mu\text{V}/\text{K}$  at  $T_{\text{min}} = 8$  K, followed by a linear dependence  $S(T) \sim T$  between 4–7 K. Although the absolute value of TEP of  $\text{CeNi}_4\text{Si}$  is small, it shows a similar structure as many other Ce-based compounds do, for instance, antiferromagnets like  $\text{CePdSn}$  [2] and  $\text{CeAl}_2$  [3].

The TEP of  $\text{YbNi}_4\text{Si}$  below 20 K exhibits rather a linear dependence on temperature (Fig. 1b). The linearity of the  $S(T)$  curve of  $\text{YbNi}_4\text{Si}$  for  $T > 200$  K is doubtlessly related to the diffusion thermopower and the phonon drag certainly

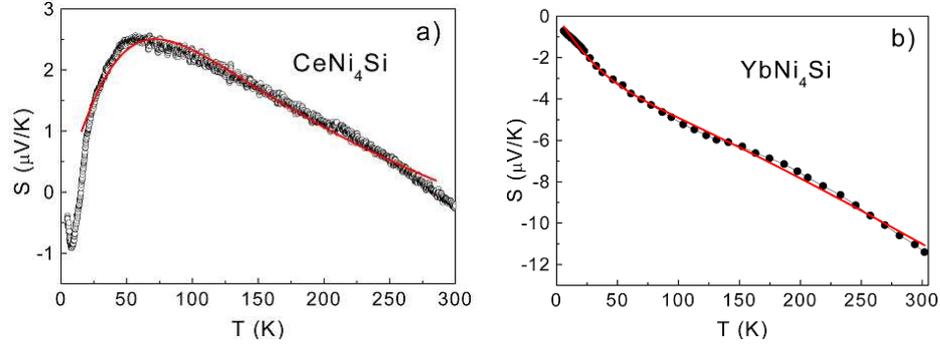


Fig. 1. Temperature dependence of the thermoelectric power of: (a)  $CeNi_4Si$ , (b)  $YbNi_4Si$ . The solid lines are fits of the experimental data with the modified Mott's equation.

plays a minor role. Besides, as is typical of Ce- and Yb-based compounds TEP is mostly positive for the former and negative for the later [4]. It is a direct consequence of the electronic structure (Yb being a hole counterpart of Ce). One can notice that TEP in the temperature range 25–175 K deviates from the linear dependence. Since the studied compound is nonmagnetic, we can neglect the magnon drag.

Ce compounds in the valence fluctuation regime exhibit a large maximum in  $S(T)$ . To explain this observation, Gottwick et al. [5] proposed a phenomenological model based on the two-band conductor model, which was originally proposed for 3D transition metals and alloys. The conduction electrons are assumed to be scattered by a  $4f$  quasiparticle band of a Lorentzian form. Modifying Mott's equation, the TEP is expressed as

$$S_B(T) = \frac{AT}{B^2 + T^2} \quad (1)$$

with

$$A = \frac{2\Delta}{|e|} \quad \text{and} \quad B^2 = 3 \frac{\Delta^2 + \Gamma^2}{\pi^2 k_B^2}. \quad (2)$$

The parameter  $\Delta = E_0 - E_F$  is a measure of the position of the DOS peak in respect of the Fermi level and  $\Gamma$  is the width of the  $4f$  band.

The fits presented in Fig. 1a include additionally the linear term due to the diffusion thermopower, i.e.,  $S(T) = S_B(T) + cT$ . The parameters have the values:  $\Delta = 0.23$  meV and  $\Gamma = 24$  meV, with  $c = -0.0045$  for  $CeNi_4Si$ . It is noticed that this model does not describe well the low-temperature part of  $S(T)$ . The deviation between the experimental and theoretical data is presumably due to the presence of the antiferromagnetic correlations, which usually emerge as a minimum in  $S(T)$  curve. As shown in Fig. 1b, the above-mentioned model gives a satisfactory fit to the data of the  $YbNi_4Si$  compound. For  $YbNi_4Si$ , the fit parameters are  $\Delta = -0.11$  meV and  $\Gamma = 10$  meV, with  $c = -0.0345$ . The

negative value for YbNi<sub>4</sub>Si indicates that DOS peak is supposedly just below the Fermi level. A tail of the 4f<sub>7/2</sub> peak of the Yb<sup>2+</sup> spin-orbit split doublet is visible in XPS [1].

#### 4. Summary

We have measured and analysed the temperature dependence of TEP for CeNi<sub>4</sub>Si and YbNi<sub>4</sub>Si compounds. This dependence for Ce and Yb compounds with the valence fluctuation is explained by the two-band conductor model assuming a single Lorentzian 4f band. The width of the narrow band  $\Gamma$  is of the same order of magnitude as that of the Ce and Yb valence fluctuation compounds. Nevertheless, a smaller value of  $\Gamma$  in the case of YbNi<sub>4</sub>Si may be responsible for a larger TEP in comparison with CeNi<sub>4</sub>Si.

We have found that the position of the DOS in the investigated Ce and Yb-based compounds lies in opposite side in respect of the Fermi level.

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